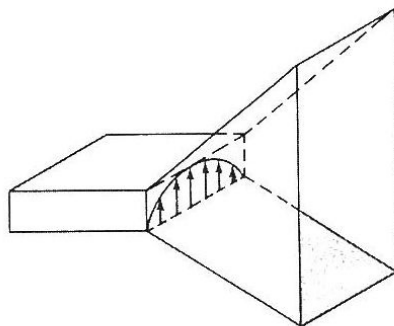


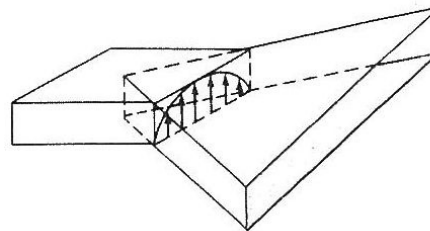


Chapter 13 : Horn Antenna

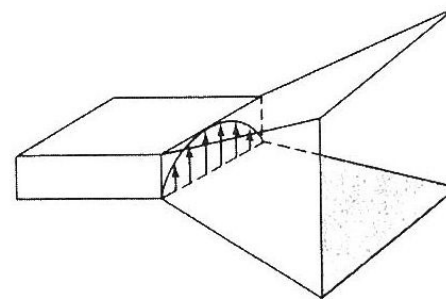
- **E-plane sectoral horn (E-plane horn)**
- **H-plane sectoral horn (H-plane horn)**
- **Pyramidal horn**
- **Corrugated horn**



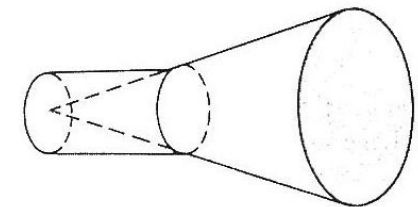
(a) E-plane



(b) H-plane



(c) Pyramidal



(d) Conical



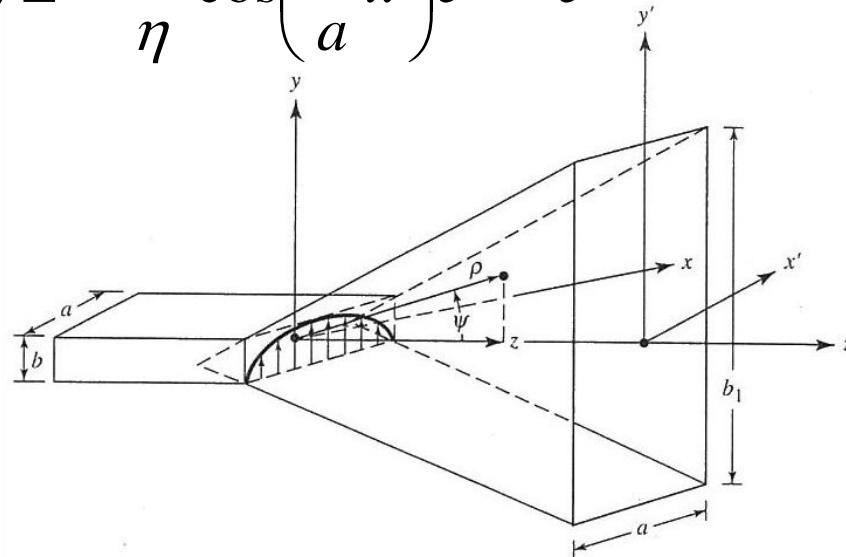
E-plane horn

The aperture fields are given by (for $-a/2 \leq x' \leq a/2$; $-b_1/2 \leq y' \leq b_1/2$)

$$E'_y(x', y') \cong E_1 \cos\left(\frac{\pi}{a} x'\right) e^{-jk\rho_1} e^{-jky'^2/2\rho_1}; E'_z = E'_y = H'_y = 0;$$

$$H'_z(x', y') \cong jE_1 \frac{\pi}{ka\eta} \sin\left(\frac{\pi}{a} x'\right) e^{-jk\rho_1} e^{-jky'^2/2\rho_1};$$

$$H'_x(x', y') \cong -\frac{E_1}{\eta} \cos\left(\frac{\pi}{a} x'\right) e^{-jk\rho_1} e^{-jky'^2/2\rho_1}$$



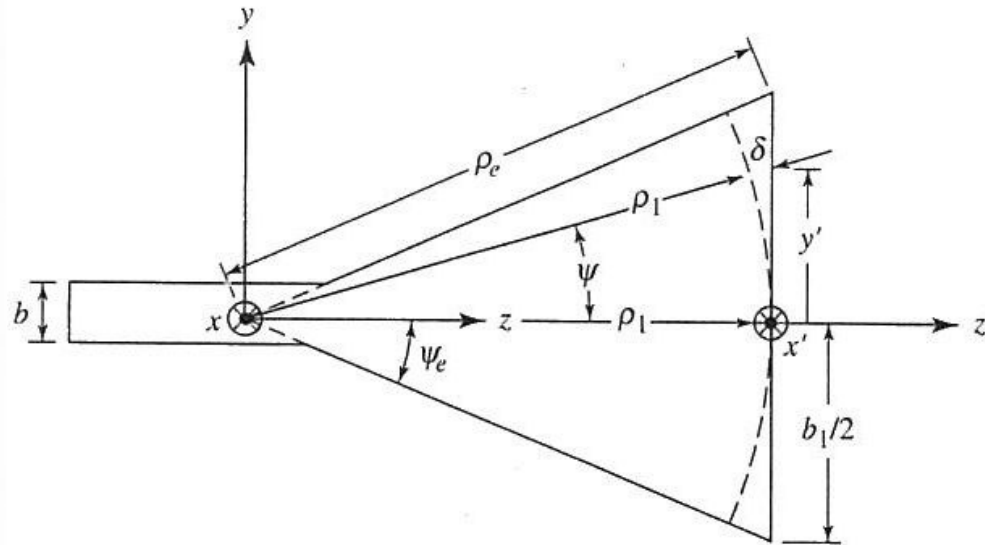


E-plane horn (2)

where $\rho_1 = \rho_e \cos \psi_e$; $2\psi_e$: flare angle

$$[\rho_1 + \delta(y')]^2 = \rho_1^2 + y'^2 \text{ or } \delta(y') = -\rho_1 + [\rho_1^2 + y'^2]^{1/2}$$

Binomial approximation:
$$\delta(y') = -\rho_1 + \rho_1 \left[1 + \frac{1}{2} \left(\frac{y'}{\rho_1} \right)^2 \right] = \frac{1}{2} \frac{y'^2}{\rho_1}$$





Design example

- Design an E-plane horn so that the maximum phase deviation at the aperture of the horn is 56.72° . The dimensions of the horn are $a=0.5\lambda$, $b_1=2.75\lambda$.

Since
$$\Delta\phi|_{\max} = k\delta(y')|_{y'=b_1/2} = \frac{k(b_1/2)^2}{2\rho_1} = 56.72\left(\frac{\pi}{180}\right)$$

or
$$\rho_1 = \left(\frac{2.75}{2}\right)^2 \frac{180}{56.72} \lambda = 6\lambda$$

Thus

$$2\psi_e = 2 \tan^{-1}\left(\frac{b_1/2}{\rho_1}\right) = 2 \tan^{-1}\left(\frac{2.75/2}{6}\right) = 25.81^\circ$$



Far-zone fields

Assume that closed surface is an infinite plane that passes through the aperture of the horn, i.e., $z' = \rho_1$ plane. On this plane the fields are assumed to be zero everywhere except at the aperture. Thus,

$$\mathbf{J}_s = \hat{n} \times \mathbf{H}_a = \hat{z} \times \mathbf{H}_a; \mathbf{M}_s = -\hat{z} \times \mathbf{E}_a$$

Hence

$$\left. \begin{aligned} J_y &= -\frac{E_1}{\eta} \cos\left(\frac{\pi}{a} x'\right) e^{-jk\rho_1} e^{-jk\delta(y')} \\ M_x &= E_1 \cos\left(\frac{\pi}{a} x'\right) e^{-jk\rho_1} e^{-jk\delta(y')} \end{aligned} \right\} \begin{aligned} -a/2 &\leq x' \leq a/2 \\ -b_1/2 &\leq y' \leq b_1/2 \end{aligned}$$

and

$$\mathbf{J}_s = \mathbf{M}_s = \mathbf{0} \quad \text{elsewhere}$$



Far-zone fields (2)

Recall that the far-zone electric field due to electric and magnetic surface currents are given by:

$$\mathbf{E}(\vec{r}) \approx -\frac{jk}{4\pi} \frac{e^{-jkr}}{r} \iint_S [\mathbf{M}_s(\vec{r}') \times \hat{r} - \hat{r} \times \hat{r} \times \eta \mathbf{J}_s(\vec{r}')] e^{jk\hat{r} \cdot \vec{r}'} ds'$$

Thus

$$\mathbf{E}(\vec{r}) \approx -\frac{jk}{4\pi} \frac{e^{-jkr}}{r} \iint_S [\hat{x} M_x \times \hat{r} - \hat{r} \times \hat{r} \times \hat{y} \eta J_y] e^{jk\hat{r} \cdot \vec{r}'} ds'$$

Using

$$\hat{r} \cdot \vec{r}' = r' \cos \zeta = (x' \cos \phi + y' \sin \phi) \sin \theta + \rho_1 \cos \theta$$

$$\hat{x} \times \hat{r} = -\hat{\phi} \cos \theta \cos \phi - \hat{\theta} \sin \phi$$

$$\hat{r} \times \hat{r} \times \hat{y} = -\hat{\theta} \cos \theta \sin \phi - \hat{\phi} \cos \phi$$



Far-zone fields (3)

The far-zone electric field due to the horn antenna can be written as:

$$\mathbf{E}(\vec{r}) \approx -\frac{jk}{4\pi} \frac{e^{-jkr}}{r} \{\mathbf{L} + \mathbf{N}\}$$

where

$$\mathbf{L} = (-\hat{\phi} \cos \theta \cos \phi - \hat{\theta} \sin \phi) e^{jk\rho_1 \cos \theta} \times \int_{-a/2}^{a/2} \int_{-b_1/2}^{b_1/2} M_x(x', y') e^{jk \sin \theta (x' \cos \phi + y' \sin \phi)} dy' dx'$$

and

$$\mathbf{N} = (\hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi) e^{jk\rho_1 \cos \theta} \times \eta \int_{-a/2}^{a/2} \int_{-b_1/2}^{b_1/2} J_y(x', y') e^{jk \sin \theta (x' \cos \phi + y' \sin \phi)} dy' dx'$$



Far-zone fields (4)

Let

$$I_1 = \int_{-a/2}^{a/2} \cos\left(\frac{\pi}{a} x'\right) e^{jkx' \sin \theta \cos \phi} dx'$$

$$I_2 = \int_{-b_1/2}^{b_1/2} e^{-jk\delta(y')} e^{jky' \sin \theta \sin \phi} dy'$$

Then **N** and **L** can be written as:

$$\mathbf{N} = -(\hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi) e^{jk\rho_1(\cos \theta - 1)} E_1 I_1 I_2$$

and

$$\mathbf{L} = -(\hat{\phi} \cos \theta \cos \phi + \hat{\theta} \sin \phi) e^{jk\rho_1(\cos \theta - 1)} E_1 I_1 I_2$$



Far-zone fields (5)

It can be shown that

$$I_1 = -\frac{\pi a}{2} \left[\frac{\cos\left(\frac{ka}{2} \sin \theta \cos \phi\right)}{\left(\frac{ka}{2} \sin \theta \cos \phi\right)^2 - \left(\frac{\pi}{2}\right)^2} \right]$$

$$I_2 = \int_{-b_1/2}^{b_1/2} e^{-jky'^2/2\rho_1} e^{jky' \sin \theta \sin \phi} dy'$$

$$\begin{aligned} k_y &= k \sin \theta \sin \phi \\ &= \int_{-b_1/2}^{b_1/2} e^{-j(ky'^2/2\rho_1 - k_y y')} dy' \end{aligned}$$

$$= e^{jk_y^2 \rho_1 / 2k} \int_{-b_1/2}^{b_1/2} e^{-j(ky' - k_y \rho_1)^2 / 2k \rho_1} dy'$$



Far-zone fields (6)

Using the following change of variable

$$\sqrt{\frac{\pi}{2}}t = \sqrt{\frac{1}{2k\rho_1}}(ky' - k_y\rho_1); dt = \sqrt{\frac{k}{\pi\rho_1}}dy'$$

yields

$$\begin{aligned} I_2 &= \sqrt{\frac{\pi\rho_1}{k}} e^{jk_y^2\rho_1/2k} \int_{t_1}^{t_2} e^{-j(\pi/2)t^2} dt \\ &= \sqrt{\frac{\pi\rho_1}{k}} e^{jk_y^2\rho_1/2k} \int_{t_1}^{t_2} \left[\cos\left(\frac{\pi}{2}t^2\right) - j \sin\left(\frac{\pi}{2}t^2\right) \right] dt \\ &= \sqrt{\frac{\pi\rho_1}{k}} e^{jk_y^2\rho_1/2k} \{ [C(t_2) - C(t_1)] - j[S(t_2) - S(t_1)] \} \end{aligned}$$



Far-zone fields (7)

where

$$t_1 = \sqrt{\frac{1}{\pi k \rho_1}} \left(-\frac{k b_1}{2} - k_y \rho_1 \right); t_2 = \sqrt{\frac{1}{\pi k \rho_1}} \left(\frac{k b_1}{2} - k_y \rho_1 \right)$$

$$\underbrace{C(x) = \int_0^x \cos\left(\frac{\pi}{2} t^2\right) dt}_{\text{cosine-Fresnel integral}}; \underbrace{S(x) = \int_0^x \sin\left(\frac{\pi}{2} t^2\right) dt}_{\text{sine-Fresnel integral}}$$

$$\mathbf{N} = (\hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi) e^{j k \rho_1 (\cos \theta - 1)} E_1 \times$$

Thus,

$$\frac{\pi a}{2} \sqrt{\frac{\pi \rho_1}{k}} e^{j k_y^2 \rho_1 / 2k} \left[\frac{\cos\left(\frac{k_x a}{2}\right)}{\left(\frac{k_x a}{2}\right)^2 - \left(\frac{\pi}{2}\right)^2} \right] F(t_1, t_2)$$



Far-zone fields (8)

and

$$\mathbf{L} = (\hat{\phi} \cos \theta \cos \phi + \hat{\theta} \sin \phi) e^{jk\rho_1(\cos \theta - 1)} E_1 \times$$
$$\frac{\pi a}{2} \sqrt{\frac{\pi \rho_1}{k}} e^{jk_y^2 \rho_1 / 2k} \left[\frac{\cos\left(\frac{k_x a}{2}\right)}{\left(\frac{k_x a}{2}\right)^2 - \left(\frac{\pi}{2}\right)^2} \right] F(t_1, t_2)$$

where

$$F(t_1, t_2) = [C(t_2) - C(t_1)] - j[S(t_2) - S(t_1)]$$

$$k_x = k \sin \theta \cos \phi; \quad k_y = k \sin \theta \sin \phi$$



Far-zone fields (9)

Therefore

$$E_{\theta} = -j\sqrt{\pi\rho_1 k a} \frac{e^{-jkr}}{8r} E_1 e^{jk\rho_1(\cos\theta-1)} \sin\phi(1+\cos\theta) \times$$

$$e^{jk_y^2 \rho_1 / 2k} \left[\frac{\cos\left(\frac{k_x a}{2}\right)}{\left(\frac{k_x a}{2}\right)^2 - \left(\frac{\pi}{2}\right)^2} \right] F(t_1, t_2)$$

$$E_{\phi} = -j\sqrt{\pi\rho_1 k a} \frac{e^{-jkr}}{8r} E_1 e^{jk\rho_1(\cos\theta-1)} \cos\phi(1+\cos\theta) \times$$

$$e^{jk_y^2 \rho_1 / 2k} \left[\frac{\cos\left(\frac{k_x a}{2}\right)}{\left(\frac{k_x a}{2}\right)^2 - \left(\frac{\pi}{2}\right)^2} \right] F(t_1, t_2)$$



E-plane and H-plane

E-Plane
($\phi = \pi/2, 3\pi/2$)
(y-z plane)

$$E_{\theta} \neq 0, H_{\phi} = \frac{E_{\theta}}{\eta};$$

$$E_{\phi} = H_{\theta} = 0$$

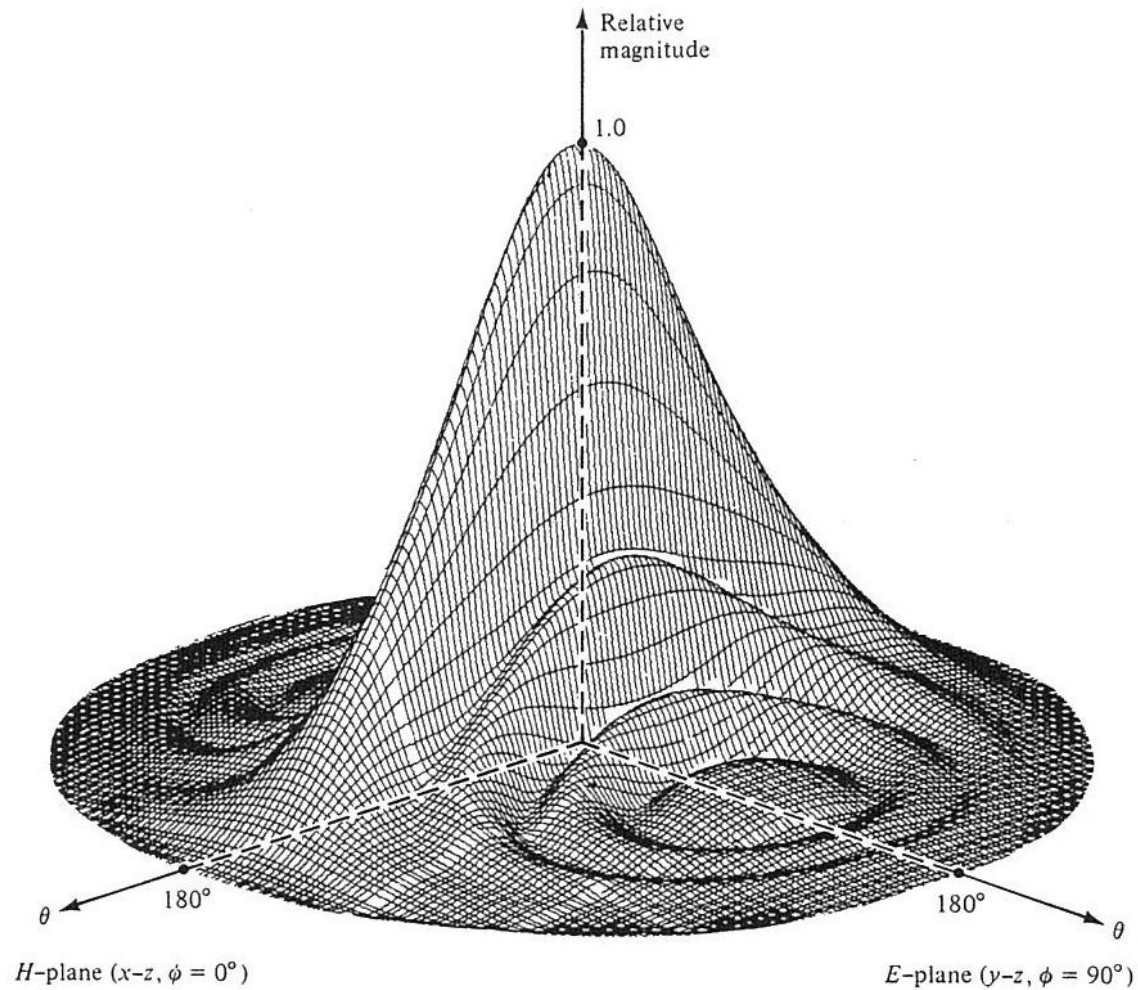
H-Plane
($\phi = 0, \pi$)
(x-z plane)

$$E_{\phi} \neq 0; H_{\theta} = -\frac{E_{\phi}}{\eta};$$

$$E_{\theta} = H_{\phi} = 0$$



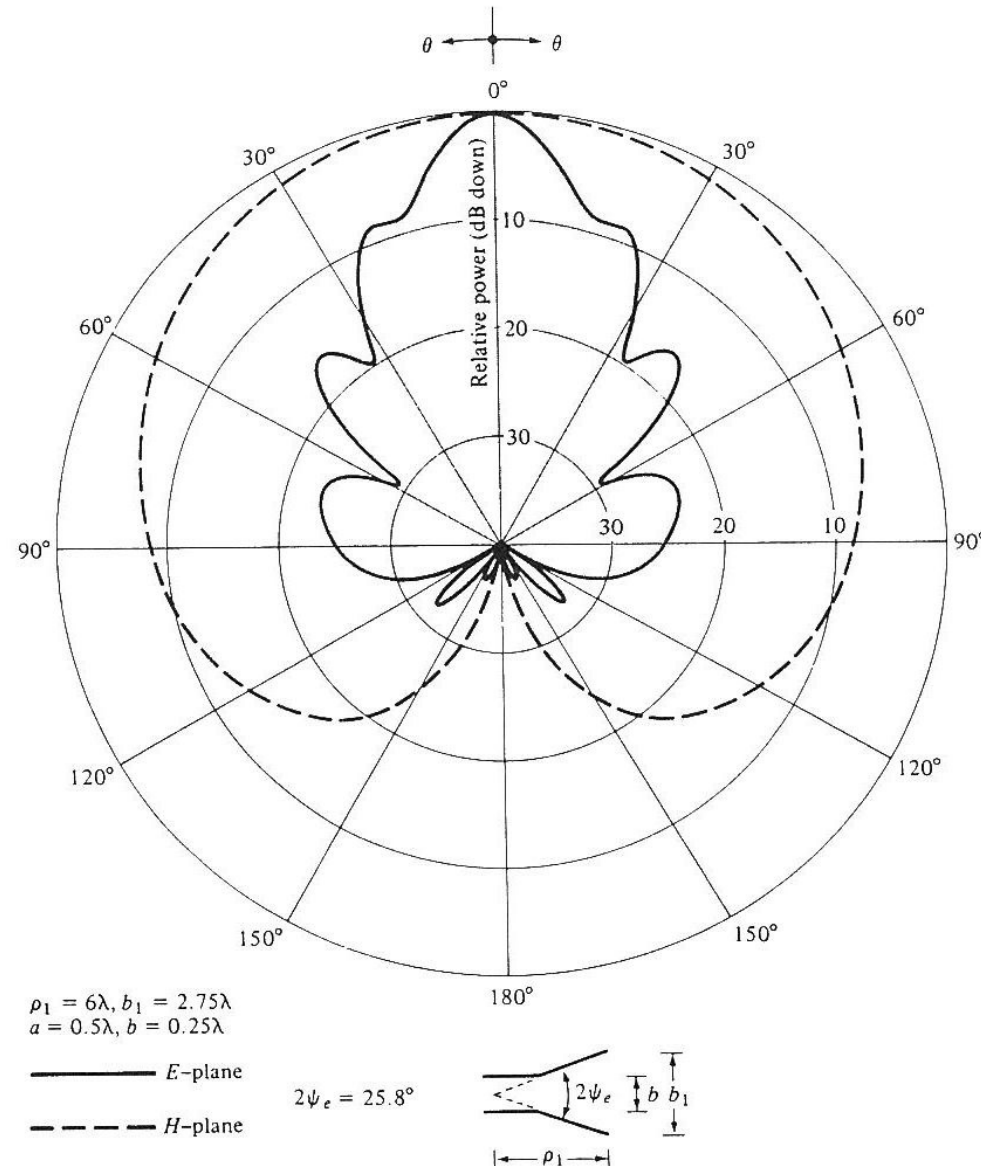
Radiation Pattern



Normalized field pattern: $\rho_1 = 6\lambda$, $a = 0.5\lambda$, $b_1 = 2.75\lambda$

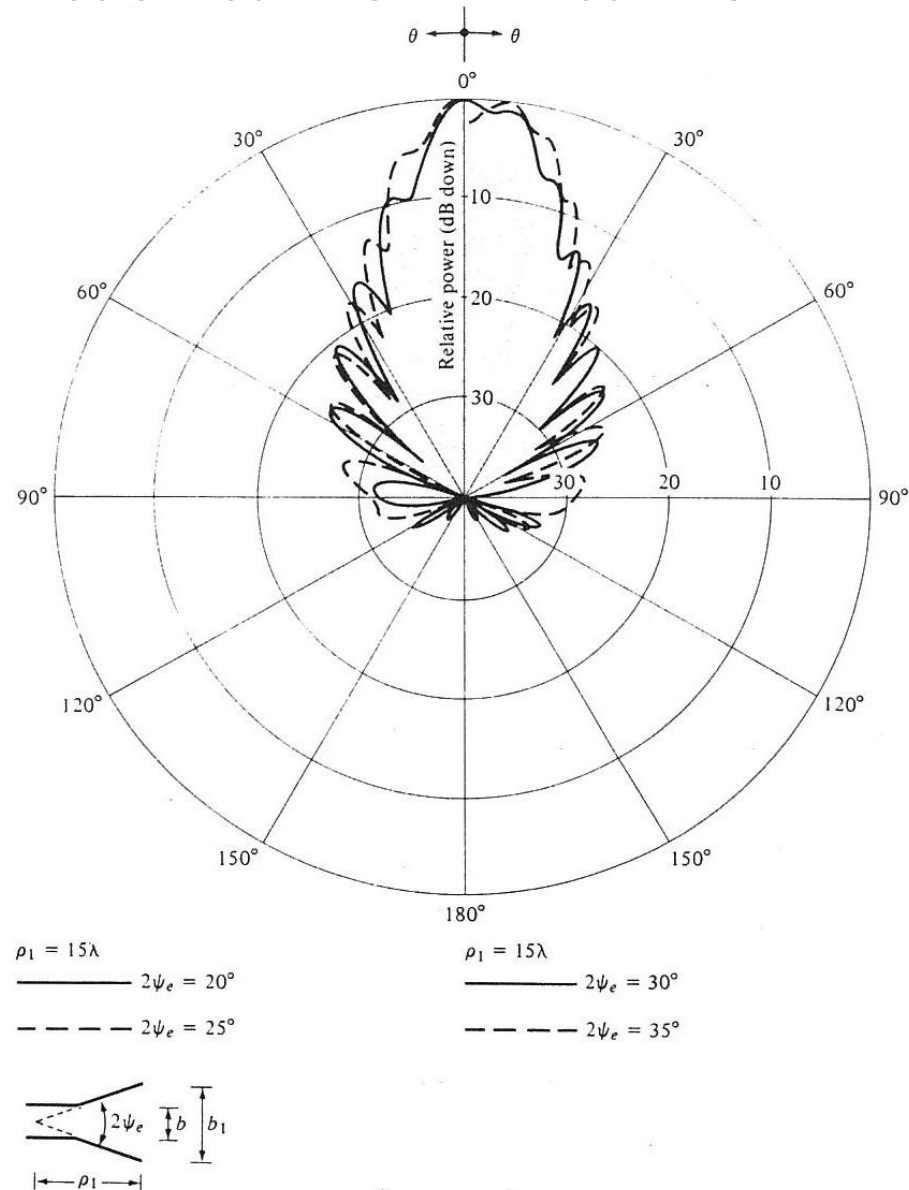


Radiation Pattern (2)





Radiation Pattern (3)





Directivity

The radiated power can be obtained from the power into the aperture. For uniform distribution, the Poynting vector is given by

$$\mathbf{W}_{av} = \frac{1}{2} \text{Re}(\mathbf{E} \times \mathbf{H}^*) = \hat{z} \frac{|E_1|^2}{2\eta} \cos^2\left(\frac{\pi}{2} x'\right)$$

for $-a/2 \leq x' \leq a/2$; $-b_1/2 \leq y' \leq b_1/2$

thus

$$P_{rad} = \int_{-a/2}^{a/2} \int_{-b_1/2}^{b_1/2} \frac{|E_1|^2}{2\eta} \cos^2\left(\frac{\pi}{2} x'\right) dy' dx' = \frac{|E_1|^2}{4\eta} ab_1$$

Recall that

$$D(\theta, \phi) = \frac{4\pi U(\theta, \phi)}{P_{rad}} \text{ and}$$

$$U(\theta, \phi) = \frac{r^2}{2\eta} (|E_\theta|^2 + |E_\phi|^2)$$



Directivity (2)

The maximum radiation intensity is given by

$$U_{\max} = U(\theta, \phi)|_{\theta=0} = \frac{r^2}{2\eta} |\mathbf{E}|_{\max}^2$$

Since $|\mathbf{E}|_{\max}^2 = |E_{\theta}|_{\max}^2 + |E_{\phi}|_{\max}^2 = \frac{4a^2 k \rho_1}{\pi^3 r^2} |E_1|^2 |F(t)|^2$

where

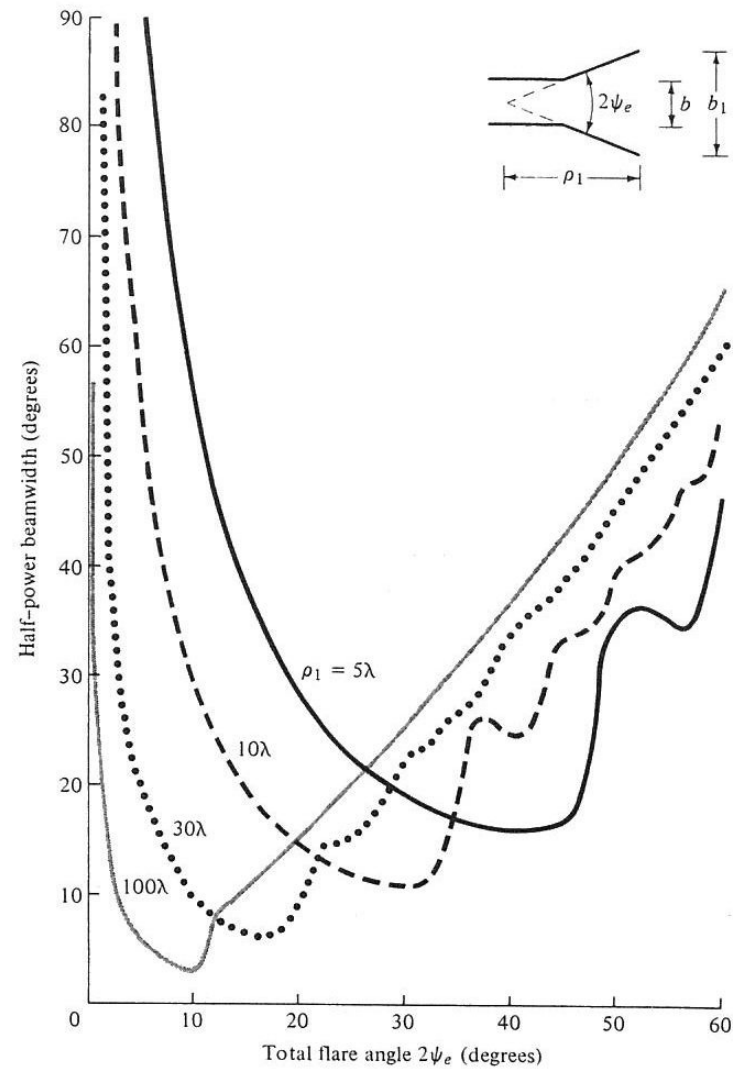
$$F(t) = C(t) - jS(t); t = \frac{b_1}{\sqrt{2\lambda\rho_1}}$$

Therefore, the maximum directivity is given by

$$\begin{aligned} D_{\max} &= \frac{4\pi U_{\max}}{P_{\text{rad}}} = \frac{64a\rho_1}{\pi\lambda b_1} |F(t)|^2 \\ &= \frac{64a\rho_1}{\pi\lambda b_1} [C^2(t) + S^2(t)] \end{aligned}$$

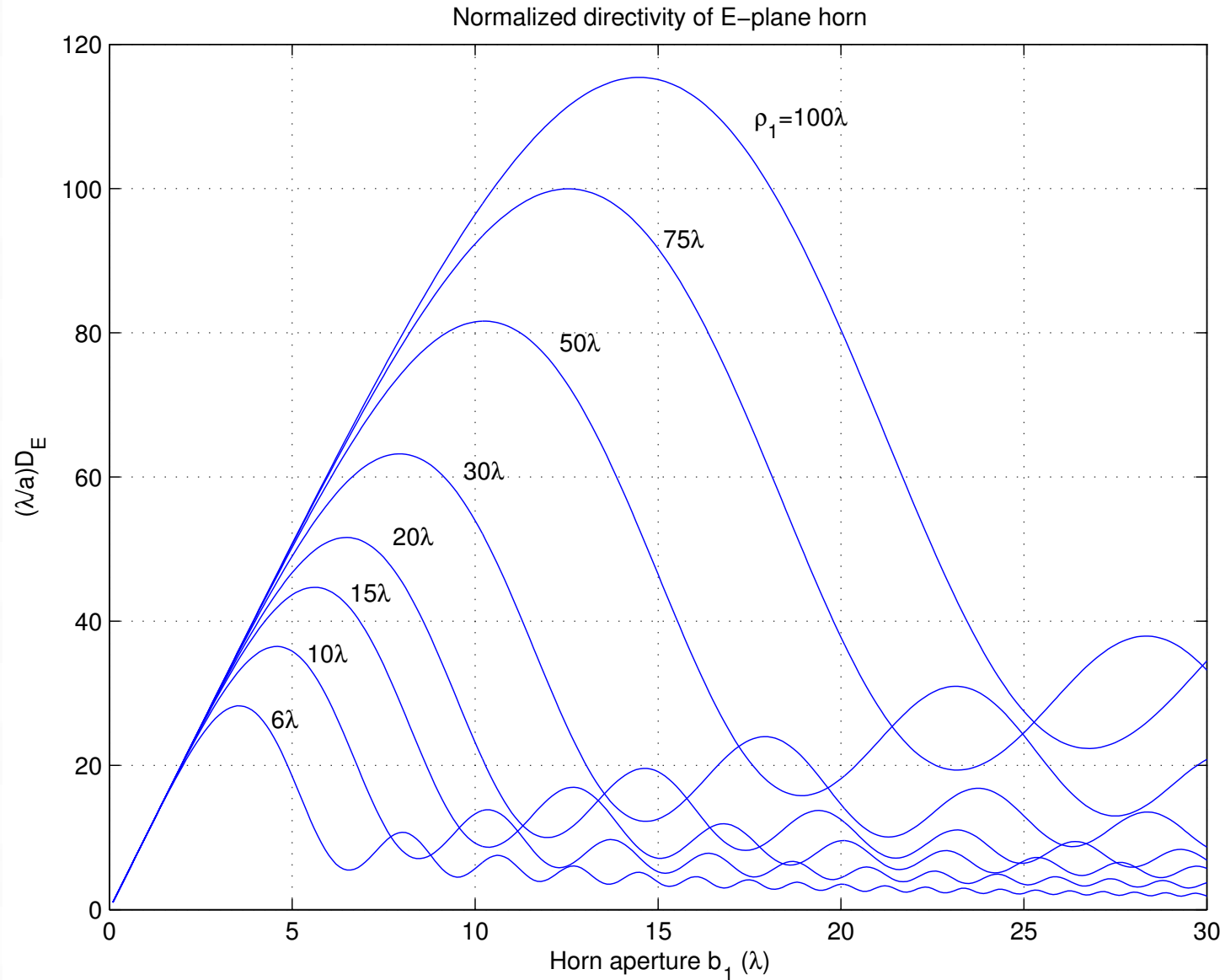


Half-power beamwidth



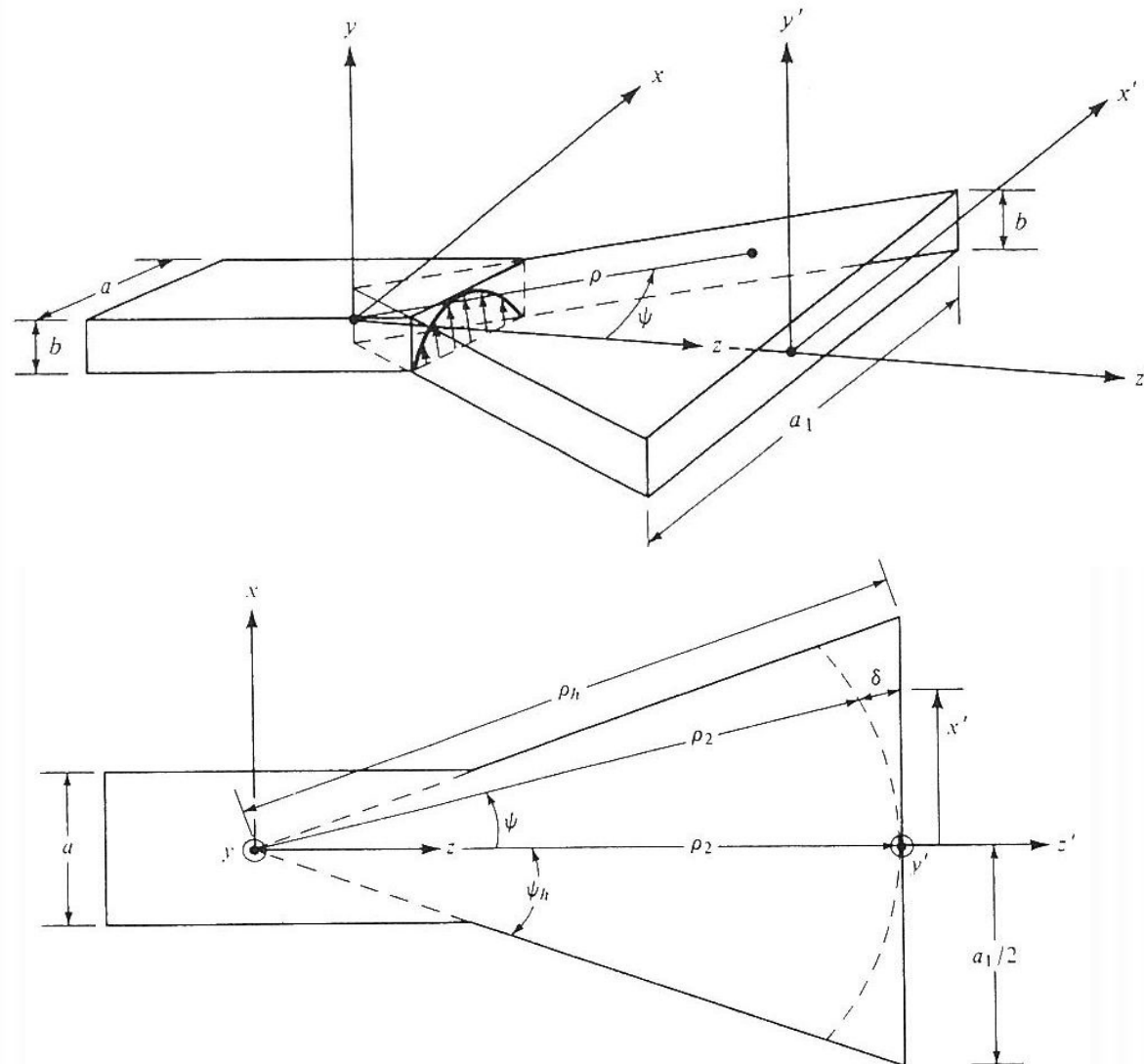


Normalized Directivity



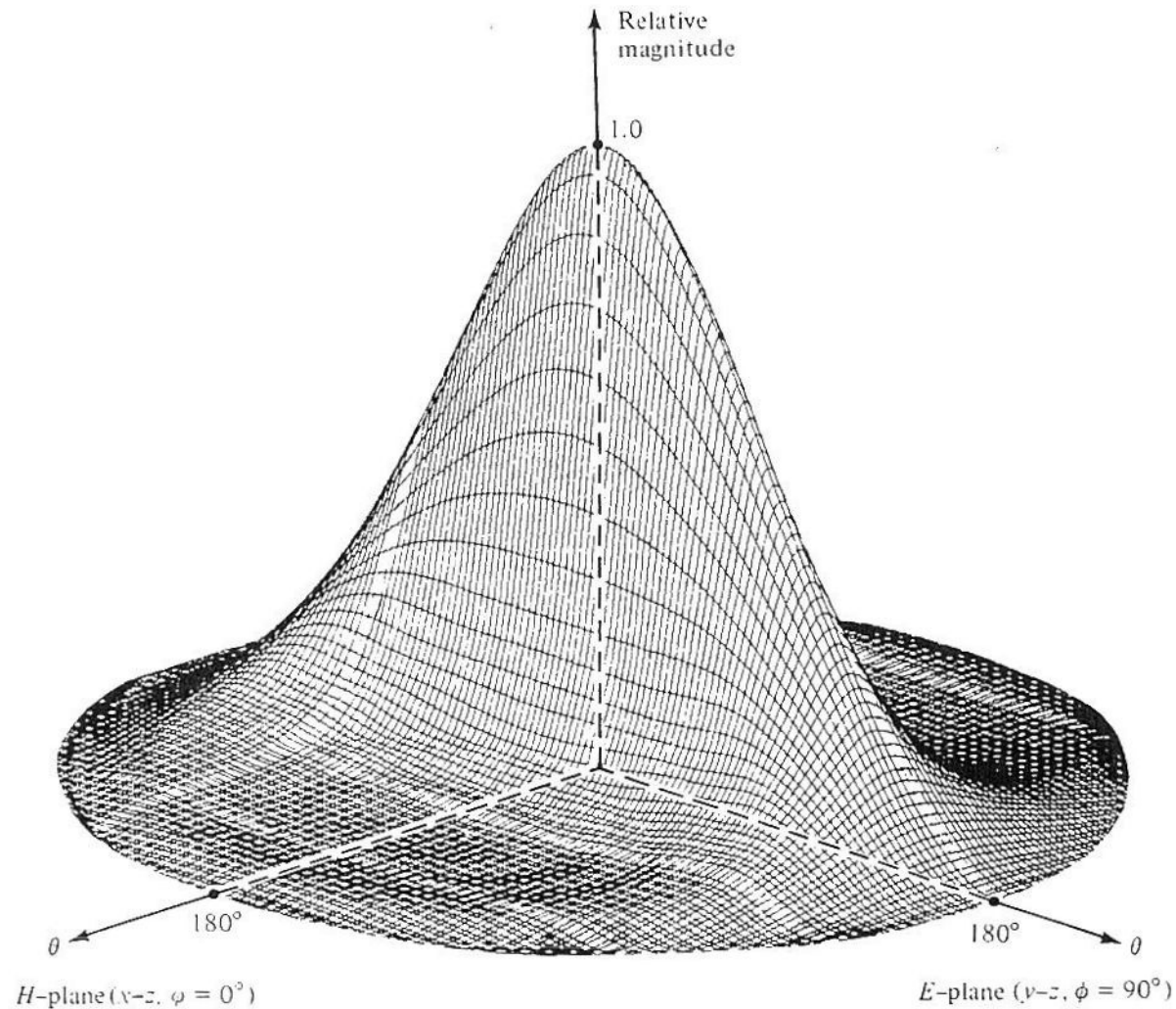


H-plane sectoral horn





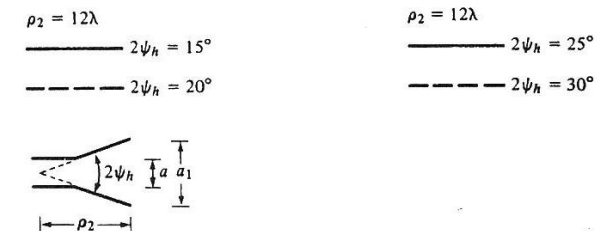
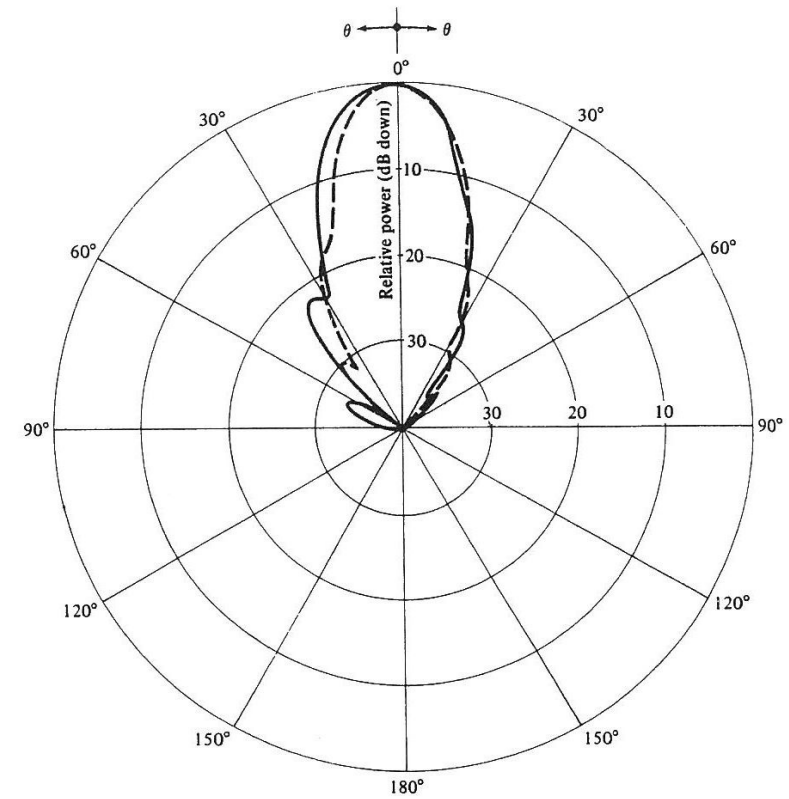
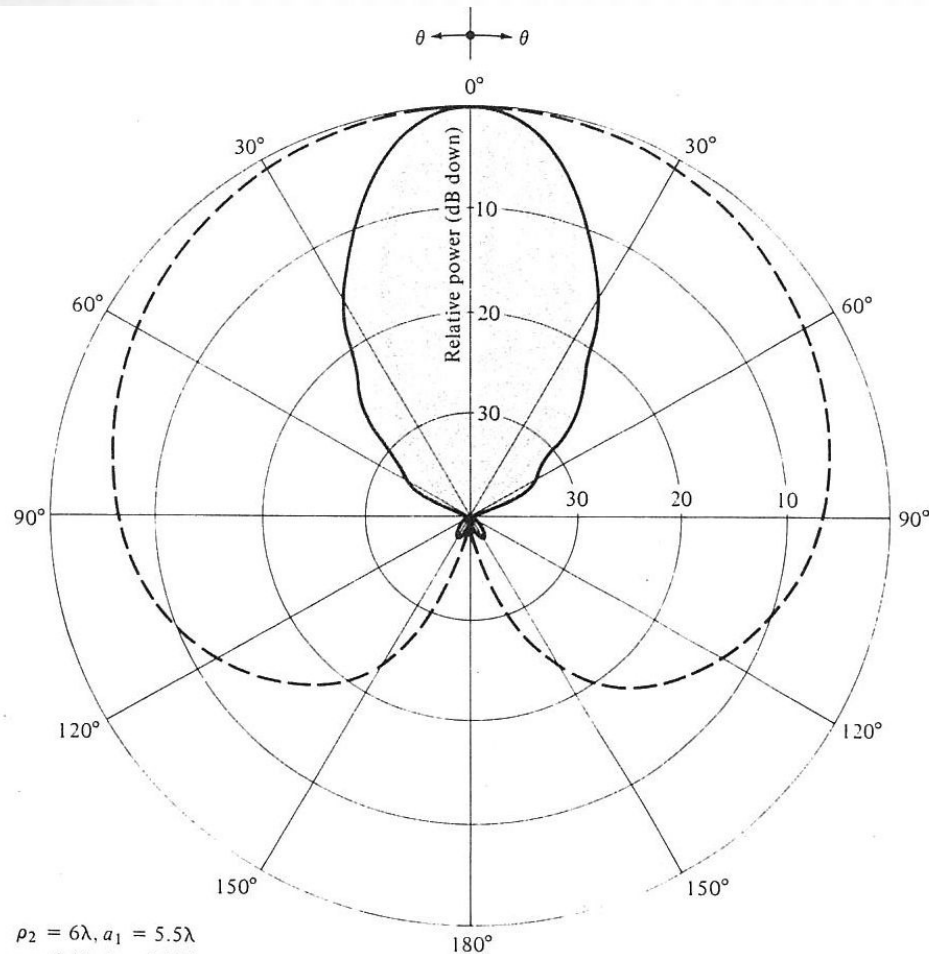
Radiation Pattern



Normalized field pattern: $\rho_2 = 6\lambda$, $a = 5.5\lambda$, $b_1 = 0.25\lambda$

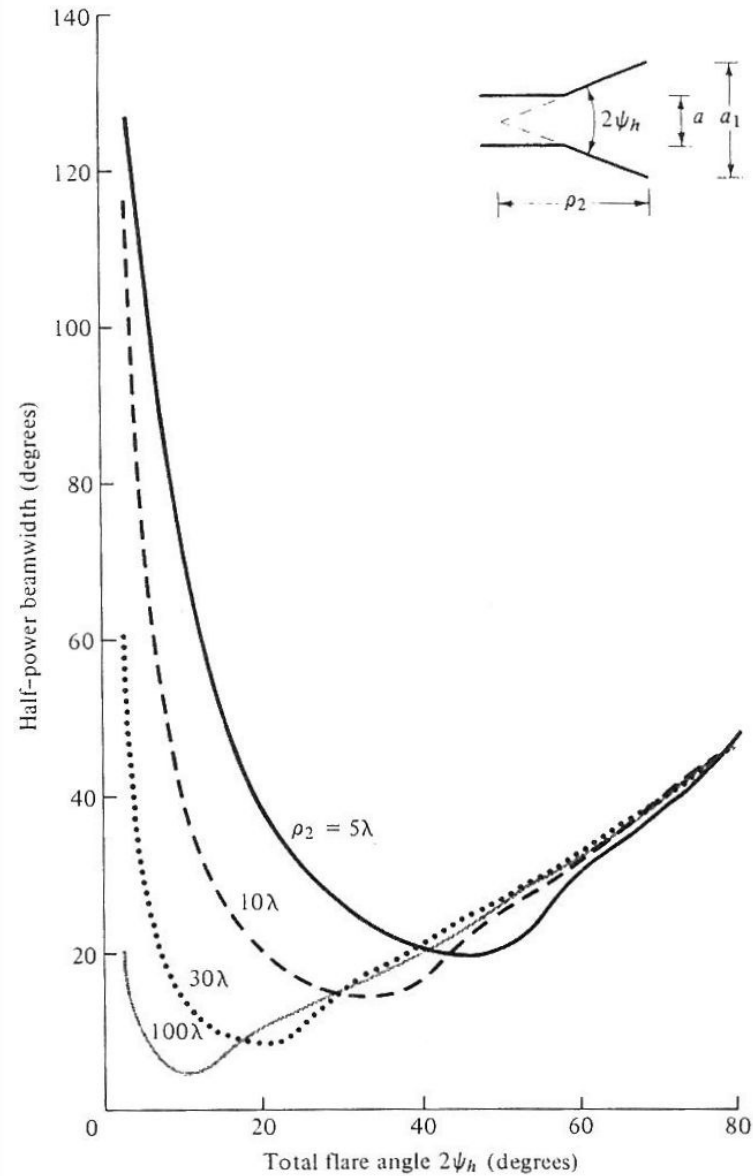


Radiation Pattern (2)



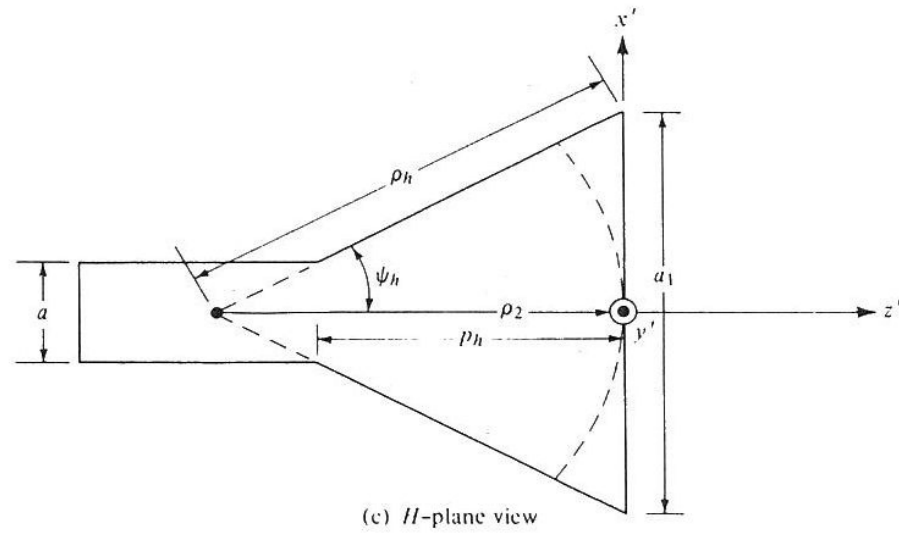
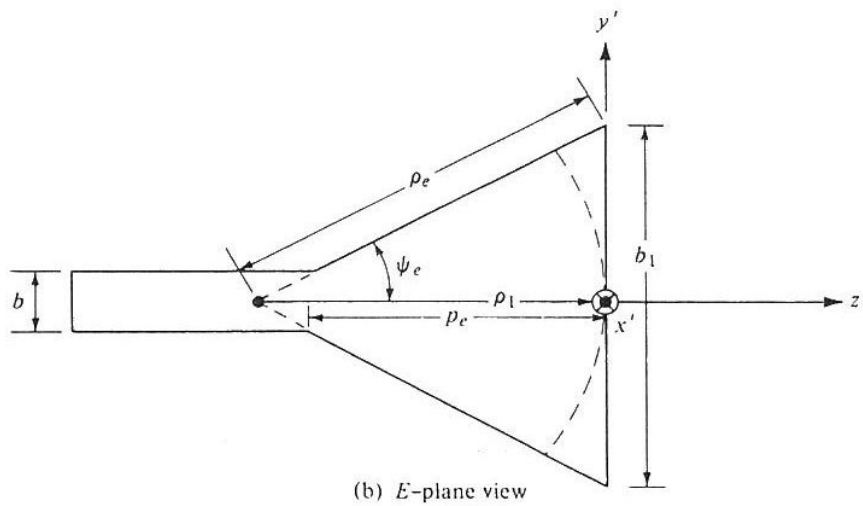
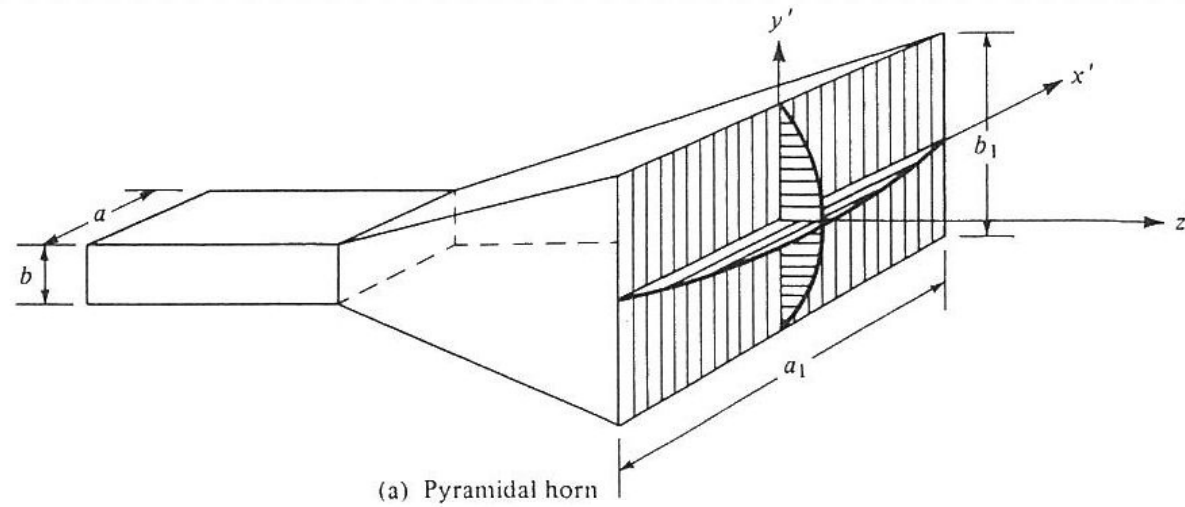


Half-power beamwidth



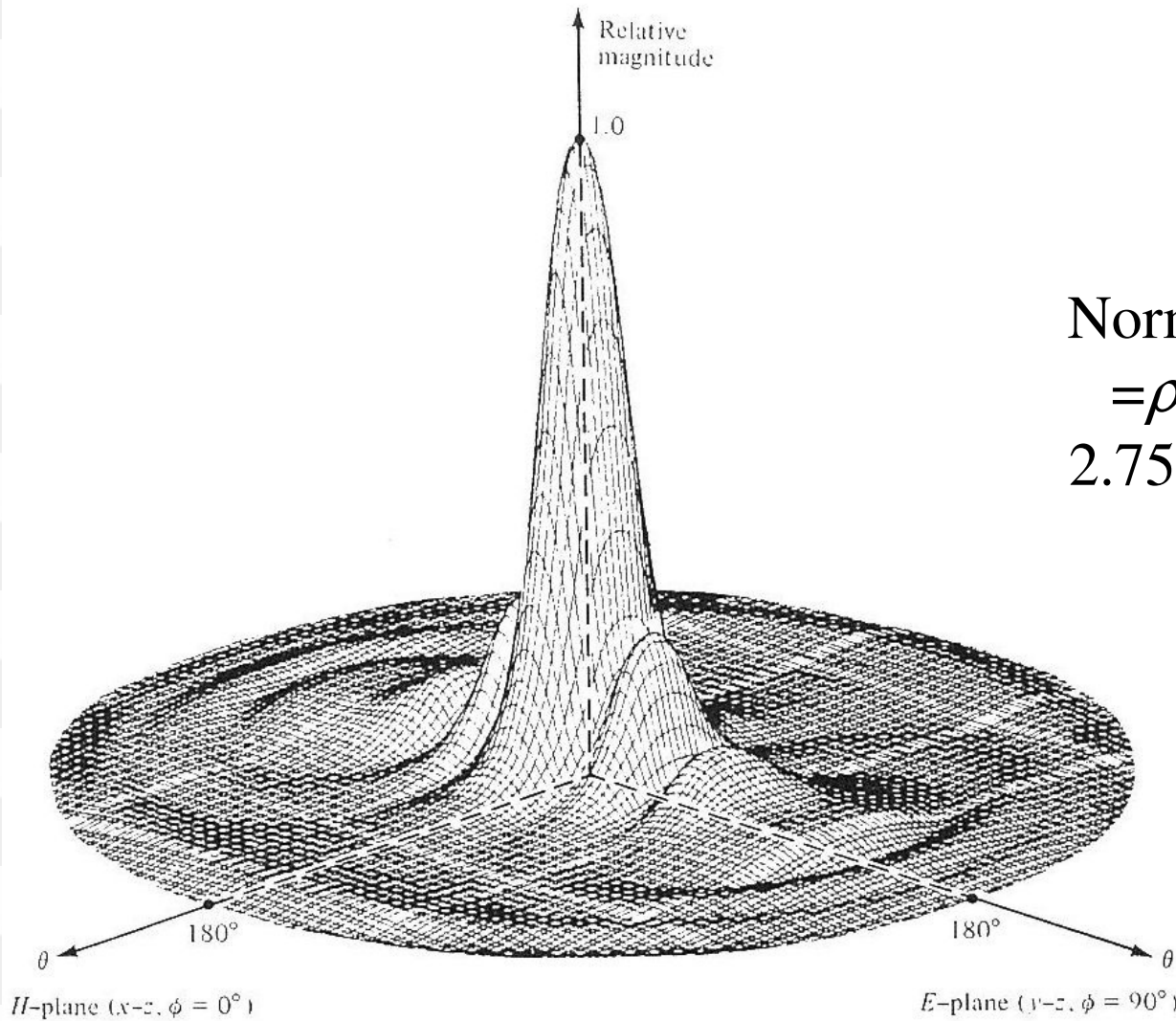


Pyramidal Horn





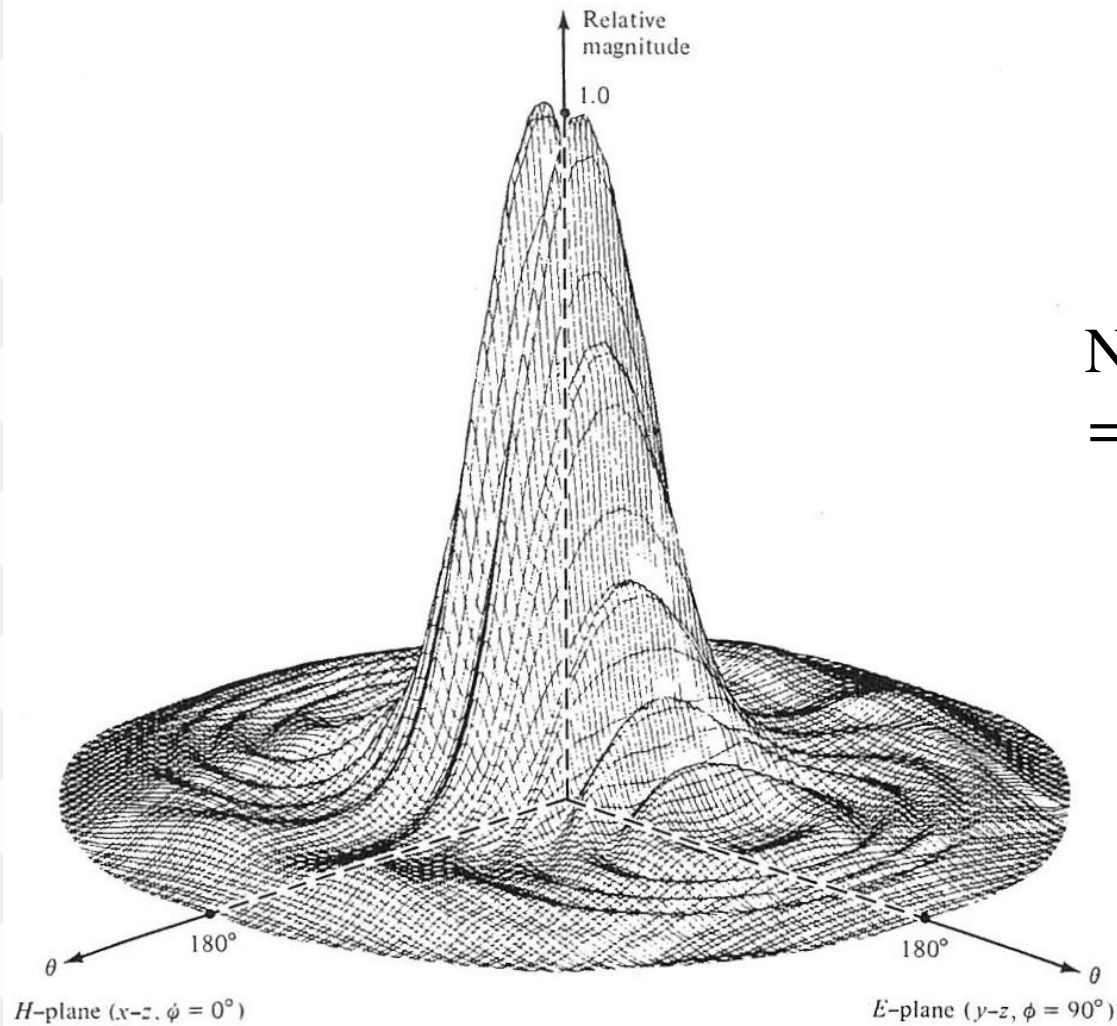
Radiation Pattern



Normalized field pattern: $\rho_1 = \rho_2 = 6\lambda$, $a_1 = 5.5\lambda$, $b_1 = 2.75\lambda$, $a_2 = 0.5\lambda$, $b_2 = 0.25\lambda$



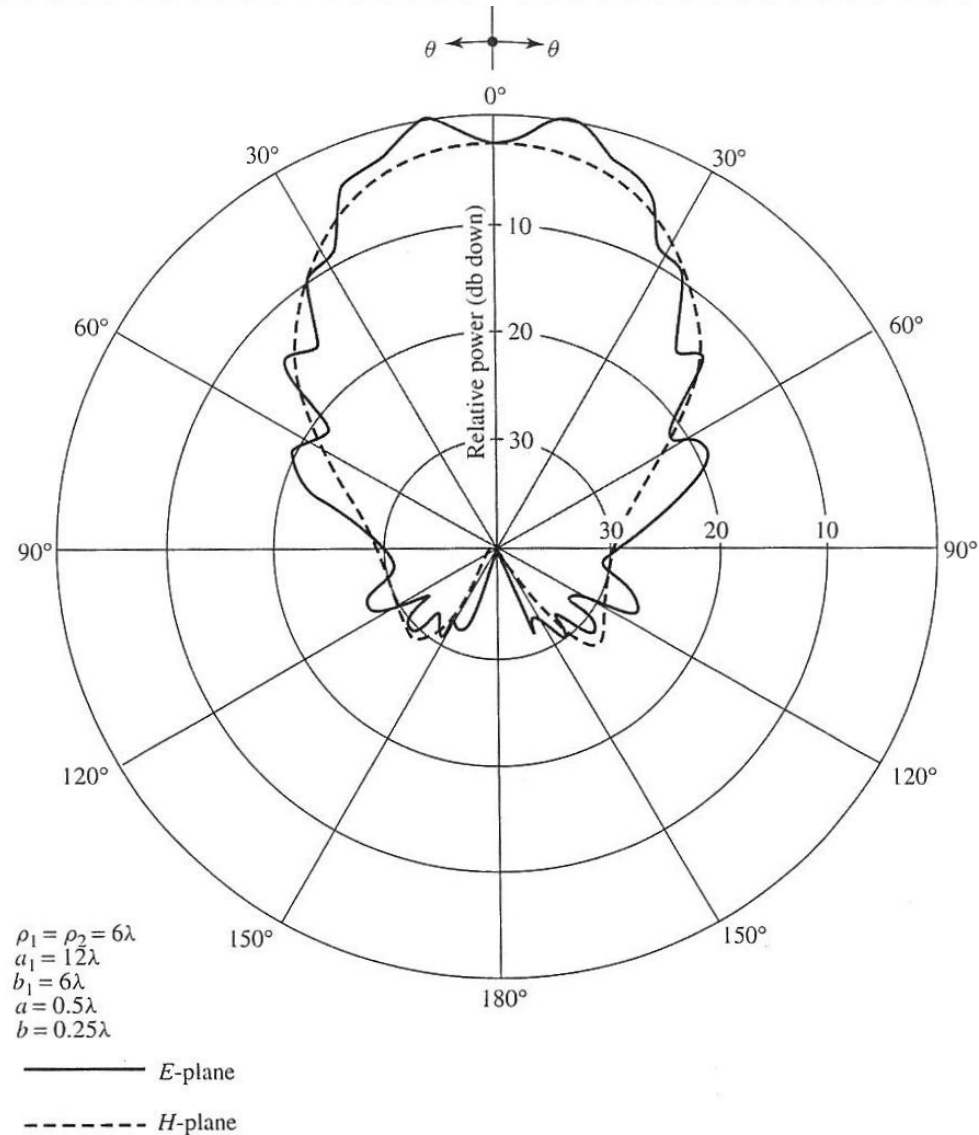
Radiation Pattern (2)



Normalized field pattern: $\rho_1 = \rho_2 = 6\lambda$, $a_1 = 12\lambda$, $b_1 = 6\lambda$,
 $a_1 = 0.5\lambda$, $b_1 = 0.25\lambda$

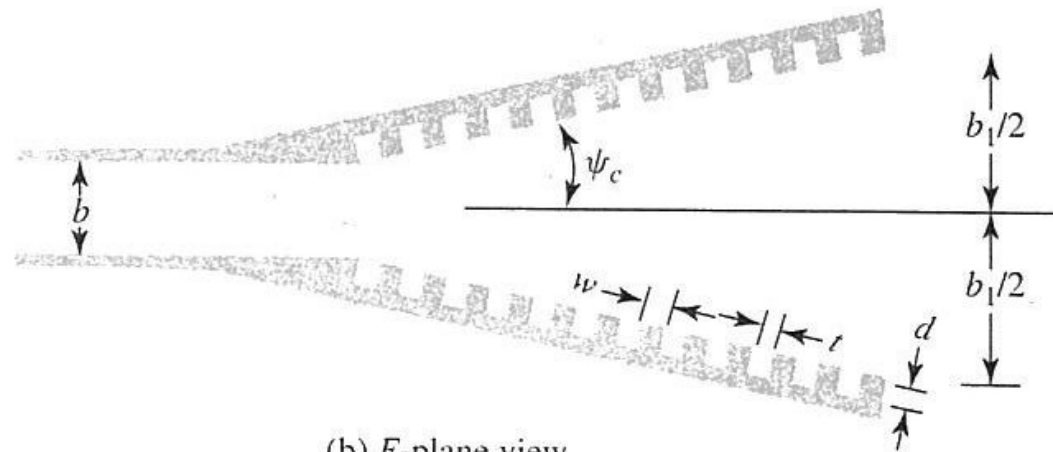
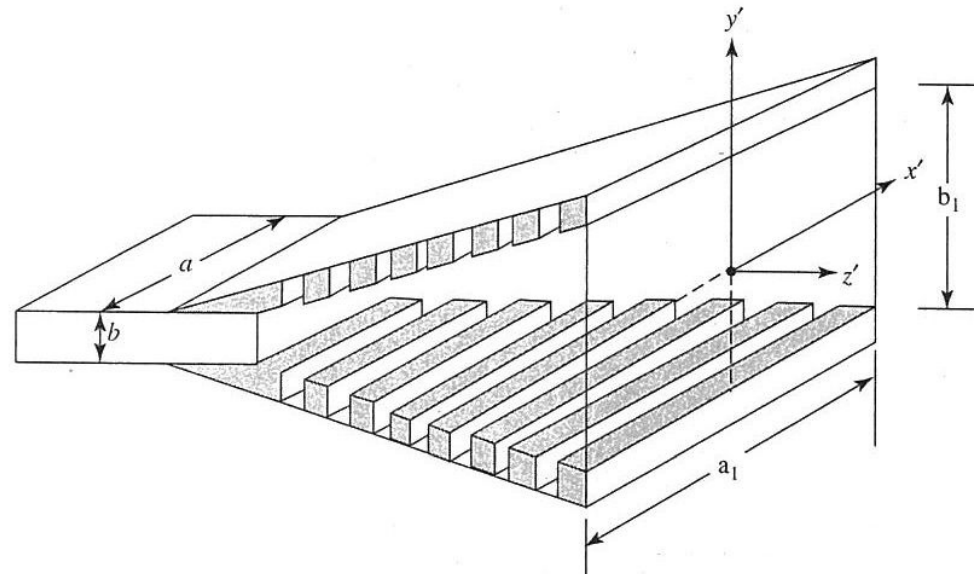


Radiation Pattern (3)





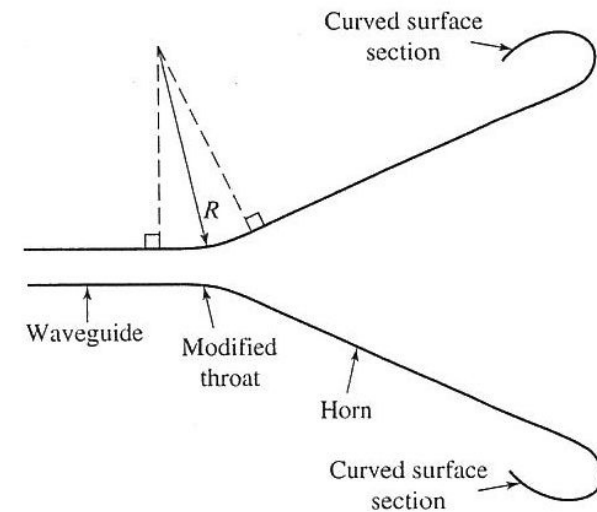
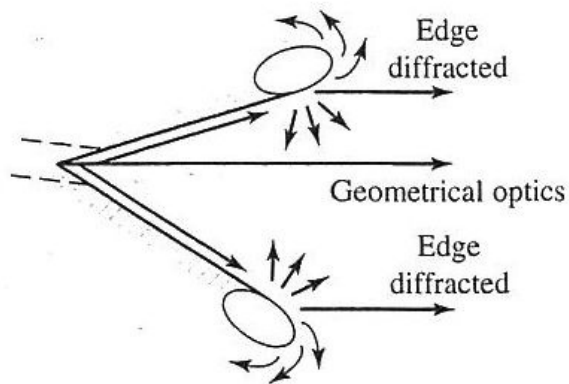
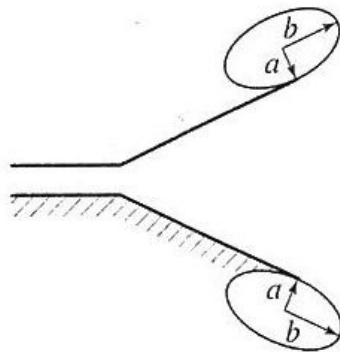
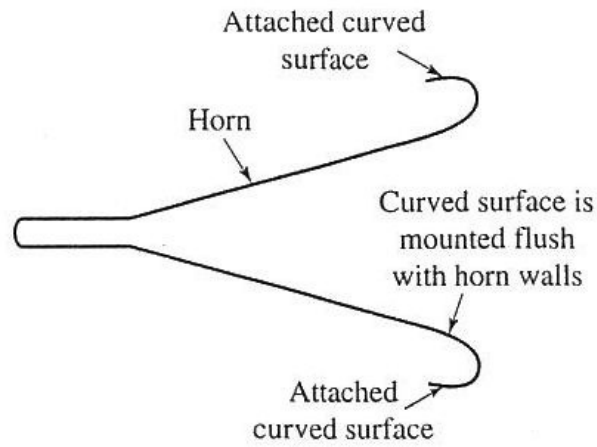
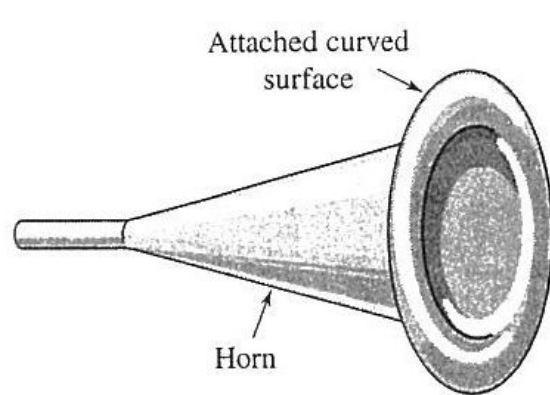
Corrugated Horns



(b) *E*-plane view



Aperture-matched horns





Quiz

- Derive the formula for the maximum directivity of E-plane horn (in slide 19).

HINT: $C(-x) = -C(x)$, $S(-x) = -S(x)$.