



Chapter 7 : Antenna Synthesis

- Continuous sources vs. Discrete sources
- Schelkunoff polynomial method
- Fourier transform method
- Woodward-Lawson method
- Triangular, cosine and cosine-squared amplitude distributions



Continuous sources

Recall the array factor

$$\mathbf{AF} = \sum_{n=1}^N a_n e^{j(n-1)\psi}; \psi = kd \cos \theta + \beta$$

If the number of elements increases in a fixed-length array, the source approaches a continuous distribution.

In the limit, the array factor becomes the space factor, i.e.,

$$\mathbf{SF} = \int_{-l/2}^{l/2} I_n(z') e^{j[kz' \cos \theta + \phi_n(z')]} dz'$$

The radiation characteristics of continuous sources can be approximated by discrete-element arrays, i.e.,

$$a_n e^{j(n-1)\beta} = I_n(z') e^{j\phi_n(z')}$$



Schelkunoff polynomial method

The array factor for an N -element, equally spaced, non-uniform amplitude, and progressive phase excitation is given by

$$\mathbf{AF} = \sum_{n=1}^N a_n e^{j(n-1)\psi}; \psi = kd \cos \theta + \beta$$

Let $z = x + jy = e^{j\psi} = e^{j(kd \cos \theta + \beta)}$

$$\mathbf{AF} = \sum_{n=1}^N a_n z^{n-1} = a_1 + a_2 z + \cdots + a_N z^{N-1}$$

which is a polynomial of degree $(N-1)$.



Schelkunoff polynomial method (2)

Thus

$$\mathbf{AF} = a_N(z - z_1)(z - z_2) \cdots (z - z_{N-1})$$

where z_1, z_2, \dots, z_{N-1} are the roots. The magnitude then becomes

$$|\mathbf{AF}| = |a_N| |z - z_1| |z - z_2| \cdots |z - z_{N-1}|$$

Note that

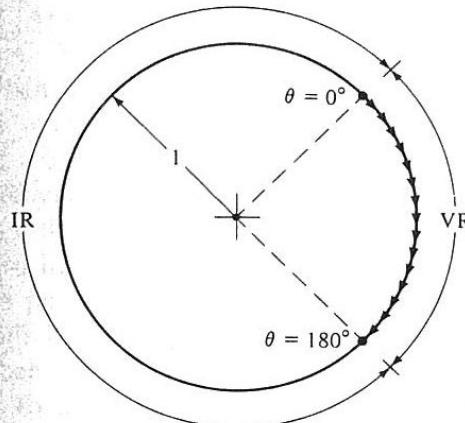
$$z = |z| e^{j\psi} = |z| \angle \psi = 1 \angle \psi$$

$$\psi = kd \cos \theta + \beta = \frac{2\pi}{\lambda} d \cos \theta + \beta$$

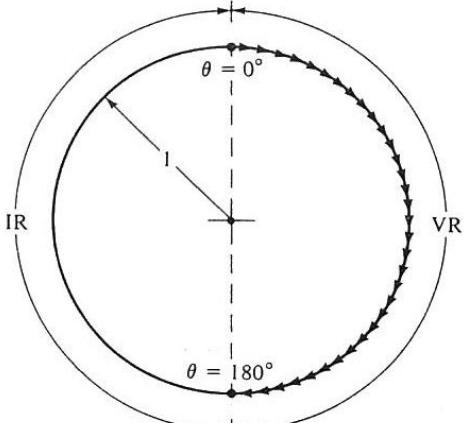
z is on a unit circle.



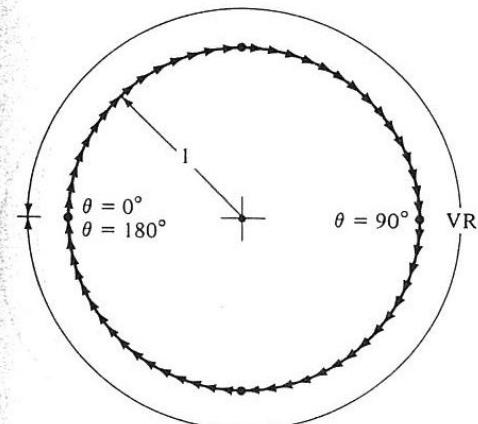
Schelkunoff polynomial method (3)



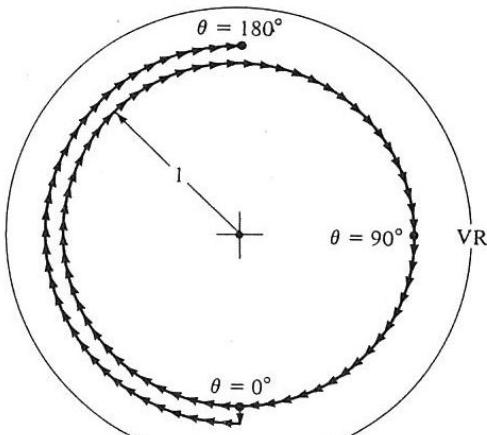
(a) $d = \lambda/8, \beta = 0$
 $\psi = \frac{\pi}{4} \cos\theta$



(b) $d = \lambda/4, \beta = 0$
 $\psi = \frac{\pi}{2} \cos\theta$



(c) $d = \lambda/2, \beta = 0$
 $\psi = \pi \cos\theta$



(d) $d = 3\lambda/4, \beta = 0$
 $\psi = \frac{3\pi}{2} \cos\theta$

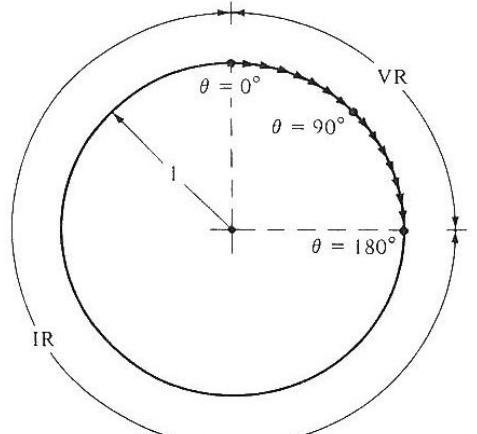
VR=Visible Region

IR=Invisible Region

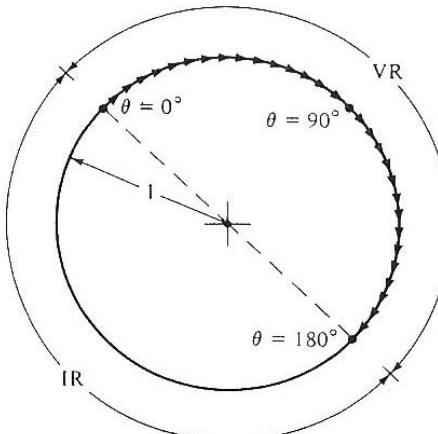
$$\beta = 0$$



Schelkunoff polynomial method (4)



(a) $d = \lambda/8, \beta = \pi/4$
 $\psi = \frac{\pi}{4} \cos\theta + \frac{\pi}{4}$

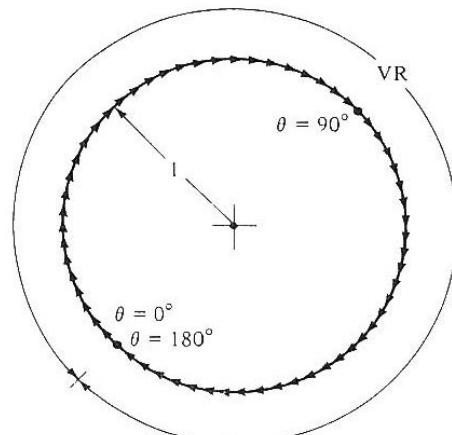


(b) $d = \lambda/4, \beta = \pi/4$
 $\psi = \frac{\pi}{2} \cos\theta + \frac{\pi}{4}$

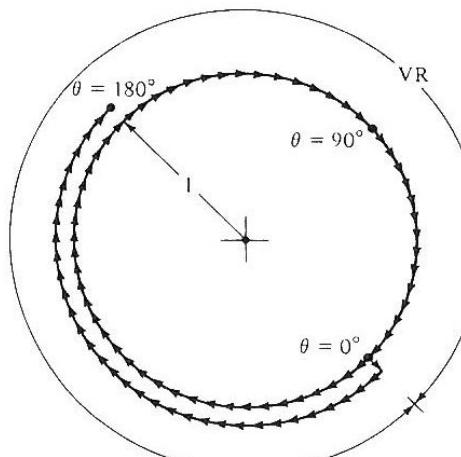
VR=Visible Region

IR=Invisible Region

$$\beta = \pi/4$$



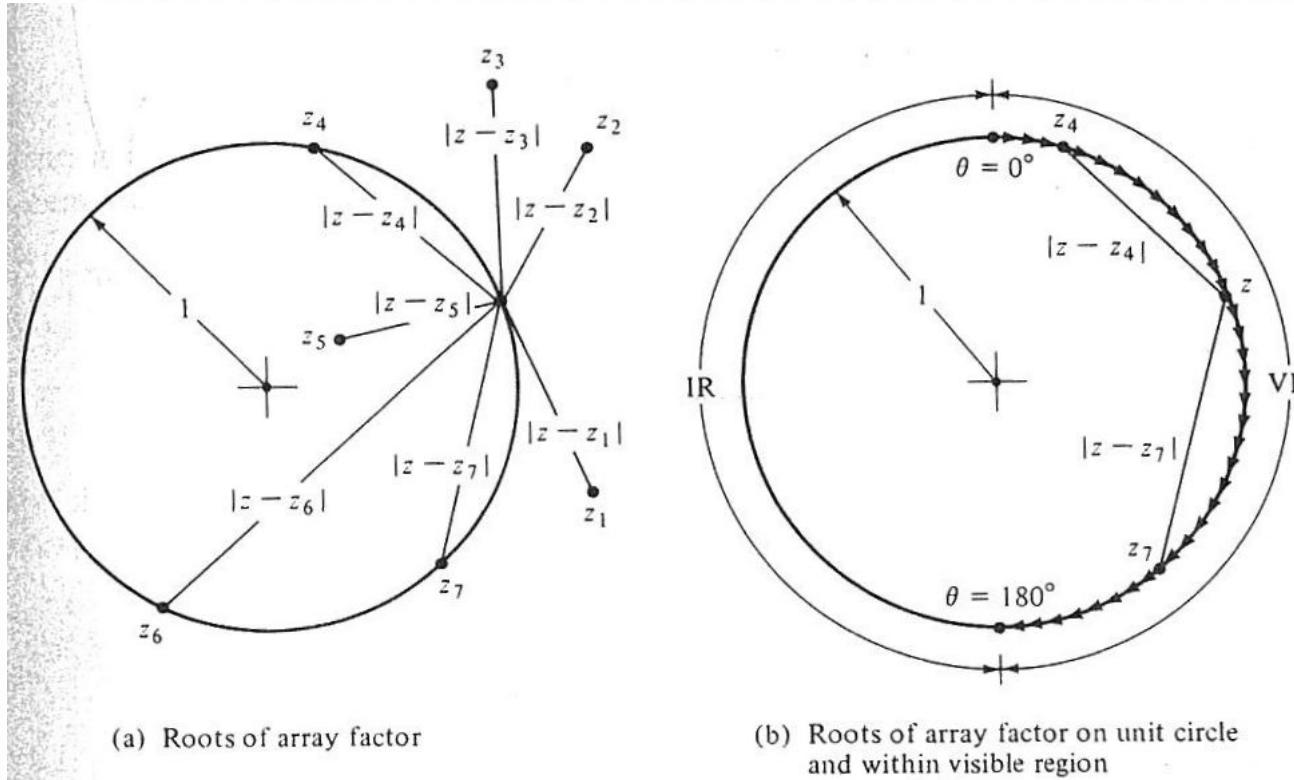
(c) $d = \lambda/2, \beta = \pi/4$
 $\psi = \pi \cos\theta + \frac{\pi}{4}$



(d) $d = 3\lambda/4, \beta = \pi/4$
 $\psi = \frac{3\pi}{2} \cos\theta + \frac{\pi}{4}$



Schelkunoff polynomial method (5)



(a) Roots of array factor

(b) Roots of array factor on unit circle
and within visible region

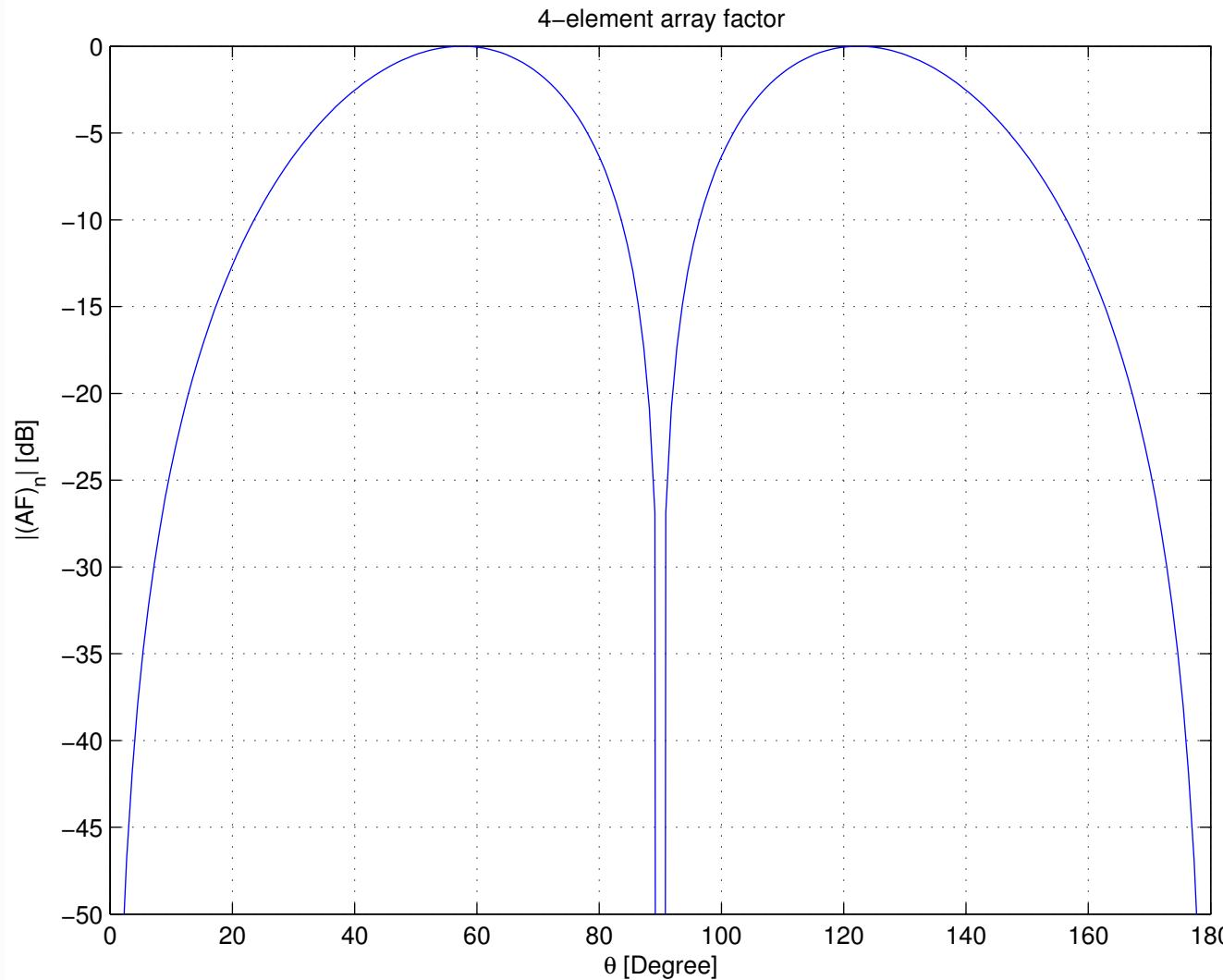


Schelkunoff polynomial method : Example

Design a linear array with a spacing between the elements of $d=\lambda/4$ such that it has zeros at $\theta=0, \pi/2, \pi$. Determine the number of elements, their excitation, and plot the derived pattern.



Schelkunoff polynomial method : Example pattern





Fourier Transform Method

The normalized space factor for a continuous line-source distribution of length l can be given by

$$\text{SF}(\theta) = \int_{-l/2}^{l/2} I(z') e^{j(k \cos \theta - k_z z')} dz' = \int_{-l/2}^{l/2} I(z') e^{j\xi z'} dz'$$

$$\xi = k \cos \theta - k_z \rightarrow \theta = \cos^{-1} \left(\frac{\xi + k_z}{k} \right)$$

where k_z is the excitation phase constant of the source. If $I(z')=I_0/l$,

$$\text{SF}(\theta) = I_0 \frac{\sin \left[\frac{kl}{2} \left(\cos \theta - \frac{k_z}{k} \right) \right]}{\frac{kl}{2} \left(\cos \theta - \frac{k_z}{k} \right)}$$



Fourier Transform Method (2)

Since the current distribution extends only over $-l/2 \leq z' \leq l/2$,

$$\mathbf{SF}(\theta) = \mathbf{SF}(\xi) = \int_{-\infty}^{\infty} I(z') e^{j\xi z'} dz'$$

The current distribution can then be given by

$$I(z') = \frac{1}{2\pi} \int_{-\infty}^{\infty} \mathbf{SF}(\xi) e^{-j\xi z'} d\xi = \frac{1}{2\pi} \int_{-\infty}^{\infty} \mathbf{SF}(\theta) e^{-j\xi z'} d\xi$$

The approximate source distribution $I_a(z')$ is given by

$$I_a(z') = \begin{cases} I(z') = \frac{1}{2\pi} \int_{-\infty}^{\infty} \mathbf{SF}(\xi) e^{-j\xi z'} d\xi & -l/2 \leq z' \leq l/2 \\ 0 & \text{elsewhere} \end{cases}$$

Thus

$$\mathbf{SF}(\theta)_a = \mathbf{SF}(\xi)_a = \int_{-l/2}^{l/2} I_a(z') e^{j\xi z'} dz'$$



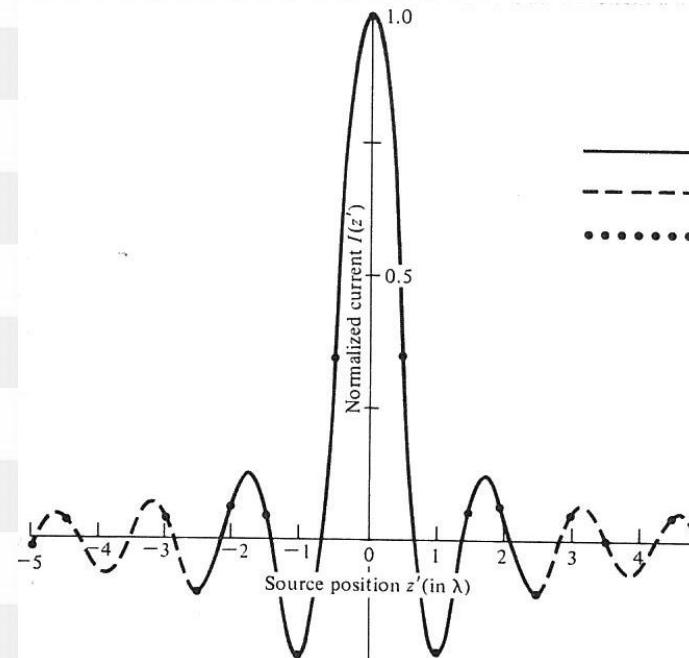
Fourier Transform Method : Example 7.2

Determine the current distribution and the approximate radiation pattern of a line source placed along the z -axis whose desired radiation pattern is symmetrical about $\theta=\pi/2$, and it is given by

$$\text{SF}(\theta) = \begin{cases} 1 & \pi/4 \leq \theta \leq 3\pi/4 \\ 0 & \text{elsewhere} \end{cases}$$

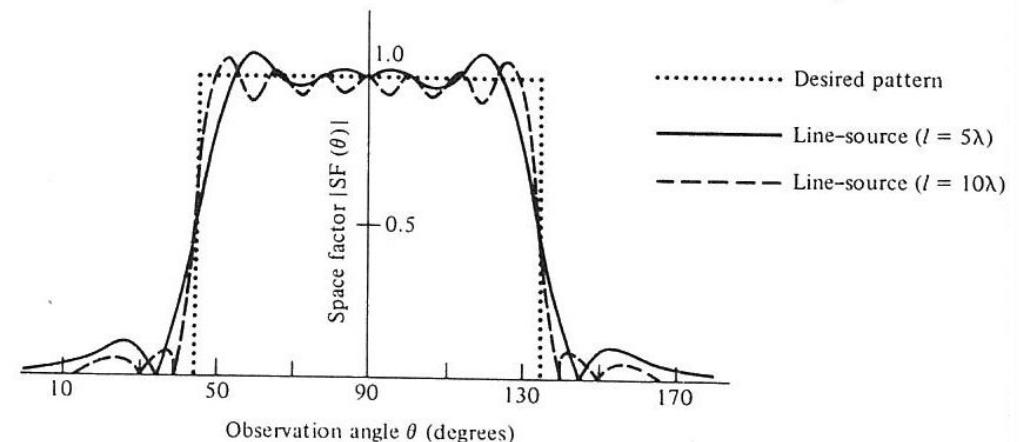


Fourier Transform Method : Example



Normalized current $I(z')$

- Line-source ($l = 5\lambda$)
- - - Line-source ($l = 10\lambda$)
- • • • • Linear array ($l = 5\lambda, 10\lambda$)



Space factor $|SF(\theta)|$

- Desired pattern
- Line-source ($l = 5\lambda$)
- - - Line-source ($l = 10\lambda$)



Fourier Transform Method : Linear Array

For an odd number of elements, the array factor is given by

$$\mathbf{AF}(\theta) = \mathbf{AF}(\psi) = \sum_{m=-M}^M a_m e^{jm\psi}$$

For an even number of elements,

$$\mathbf{AF}(\theta) = \mathbf{AF}(\psi) = \sum_{m=-M}^{-1} a_m e^{j[(2m+1)/2]\psi} + \sum_{m=1}^M a_m e^{j[(2m-1)/2]\psi}$$

where $\psi = kd \cos \theta + \beta$

Odd-number

Elements' locations

Even-number

$$z_m = \begin{cases} \frac{2m-1}{2}d & 1 \leq m \leq M \\ \frac{2m+1}{2}d & -M \leq m \leq -1 \end{cases}$$



Fourier Transform Method : Linear Array (2)

For an odd number of elements, the excitation coefficients can be obtained by

$$a_m = \frac{1}{T} \int_{-T/2}^{T/2} \mathbf{AF}(\psi) e^{-jm\psi} d\psi = \frac{1}{2\pi} \int_{-\pi}^{\pi} \mathbf{AF}(\psi) e^{-jm\psi} d\psi$$

$$-M \leq m \leq M$$

For an even number of elements,

$$a_m = \begin{cases} \frac{1}{2\pi} \int_{-\pi}^{\pi} \mathbf{AF}(\psi) e^{-j[(2m+1)/2]\psi} d\psi & -M \leq m \leq -1 \\ \frac{1}{2\pi} \int_{-\pi}^{\pi} \mathbf{AF}(\psi) e^{-j[(2m-1)/2]\psi} d\psi & 1 \leq m \leq M \end{cases}$$

where $\psi = kd \cos \theta + \beta$



Fourier Transform Method : Example

Same as Example 7.2 with $d = \lambda/2$; non-zero only

$$\pi/4 \leq \theta \leq 3\pi/4$$

thus $-\pi/\sqrt{2} \leq \psi = kd \cos \theta + \beta \leq \pi/\sqrt{2}$

therefore

$$a_m = \frac{1}{2\pi} \int_{-\pi/\sqrt{2}}^{\pi/\sqrt{2}} e^{-jm\psi} d\psi = \frac{1}{\sqrt{2}} \frac{\sin\left(\frac{m\pi}{\sqrt{2}}\right)}{m\pi}$$

$$a_0 = 1.0 \quad a_{\pm 4} = 0.0578 \quad a_{\pm 8} = -0.0496$$

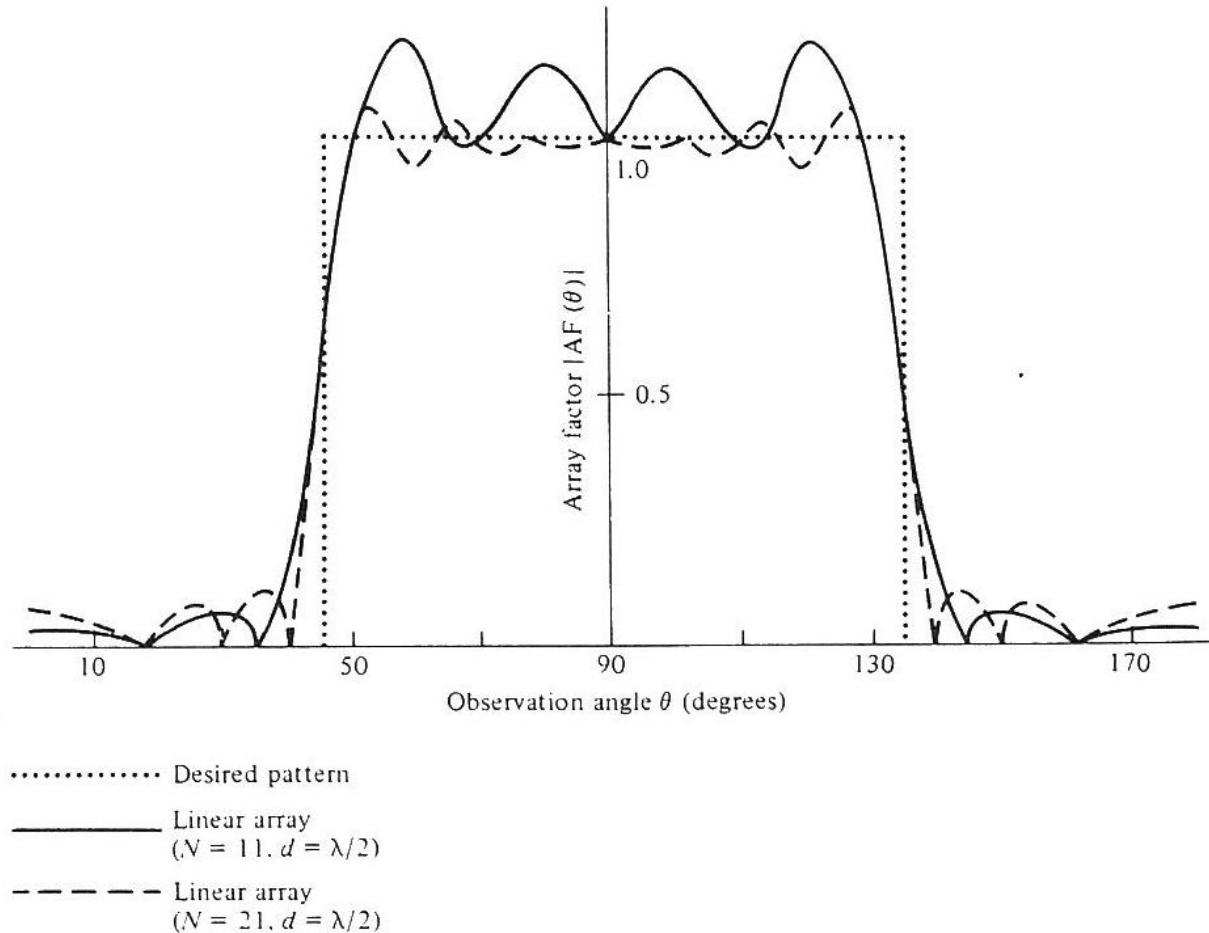
Result $a_{\pm 1} = 0.3582 \quad a_{\pm 5} = -0.0895 \quad a_{\pm 9} = 0.0455$

$$a_{\pm 2} = -0.2170 \quad a_{\pm 6} = 0.0518 \quad a_{\pm 10} = -0.010$$

$$a_{\pm 3} = 0.0588 \quad a_{\pm 7} = 0.0101$$



Fourier Transform Method : Example





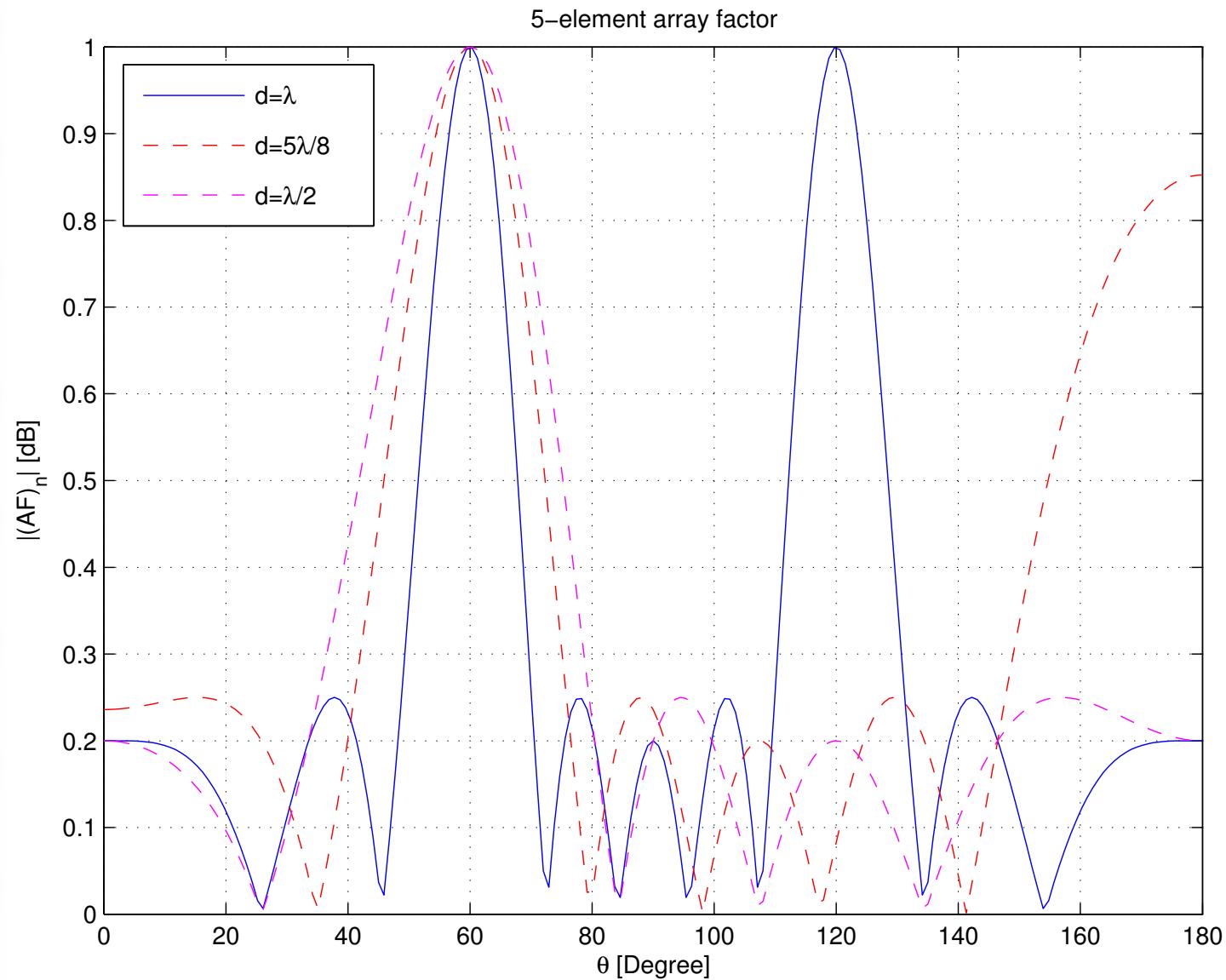
Quiz

A 5-element uniform linear array with a spacing of λ between elements is designed to scan at $\theta = \pi/3$. Assume that the array is aligned along the z-axis.

- a) Find the array factor
- b) Find the angle of the grating lobe.
- c) Find the condition such that there exists no grating lobe.



Quiz solution





Woodward-Lawson Method

- Sampling the desired pattern at various discrete locations.
- Use *composing function* of the forms:

$$b_m \sin(\psi_m)/\psi_m \text{ or } b_m \sin(N\phi_m)/N \sin \phi_m$$

as the field of each pattern sample

- The synthesized pattern is represented by a finite sum of composing functions.
- The total excitation is a sum of space harmonics.



Woodward-Lawson Method: Line-source

Let the source be represented by a sum of the following constant current source of length l .

$$i_m(z') = \frac{b_m}{l} e^{-jkz' \cos \theta_m} \quad -l/2 \leq z' \leq l/2$$

Then the current source can be given by

$$I(z') = \frac{1}{l} \sum_{m=-M}^M b_m e^{-jkz' \cos \theta_m} \quad \begin{aligned} &\text{where } m = \pm 1, \pm 2, \dots, \pm M \text{ (for } 2M \text{ even number)} \\ &m = 0, \pm 1, \pm 2, \dots, \pm M \text{ (for } 2M + 1 \text{ odd number)} \end{aligned}$$

The field pattern of each current source is given by

$$s_m(\theta) = b_m \frac{\sin\left[\frac{kl}{2}(\cos \theta - \cos \theta_m)\right]}{\frac{kl}{2}(\cos \theta - \cos \theta_m)}$$

Composing function



Woodward-Lawson Method: Line-source (2)

For an odd number samples, the total pattern becomes

$$\text{SF}(\theta) = \sum_{m=-M}^M b_m \frac{\sin\left[\frac{kl}{2}(\cos\theta - \cos\theta_m)\right]}{\frac{kl}{2}(\cos\theta - \cos\theta_m)}$$

b_m can be obtained from the value at the sample points θ_m , i.e.,

$$b_m = \text{SF}(\theta = \theta_m)_d$$

In order to satisfy the periodicity of 2π and faithfully reconstruct the desired pattern,

$$kz' \Delta \Big|_{|z'|=l} = 2\pi \Rightarrow \Delta = \frac{\lambda}{l}$$



Woodward-Lawson Method: Line-source (3)

Thus the location of each sample is given by

$$\cos \theta_m = m\Delta = m\left(\frac{\lambda}{l}\right), \quad m = 0, \pm 1, \pm 2, \dots \text{ for odd samples}$$

$$\cos \theta_m = \begin{cases} \frac{2m-1}{2}\Delta = \frac{2m-1}{2}\left(\frac{\lambda}{l}\right), & m = 1, 2, \dots \\ \frac{2m+1}{2}\Delta = \frac{2m+1}{2}\left(\frac{\lambda}{l}\right), & m = -1, -2, \dots \end{cases} \text{ for even samples}$$

Therefore, M should be the closest integer to $M=l/\lambda$.



Woodward-Lawson Method: Example

Same as Example 7.2; for $l = 5\lambda$.

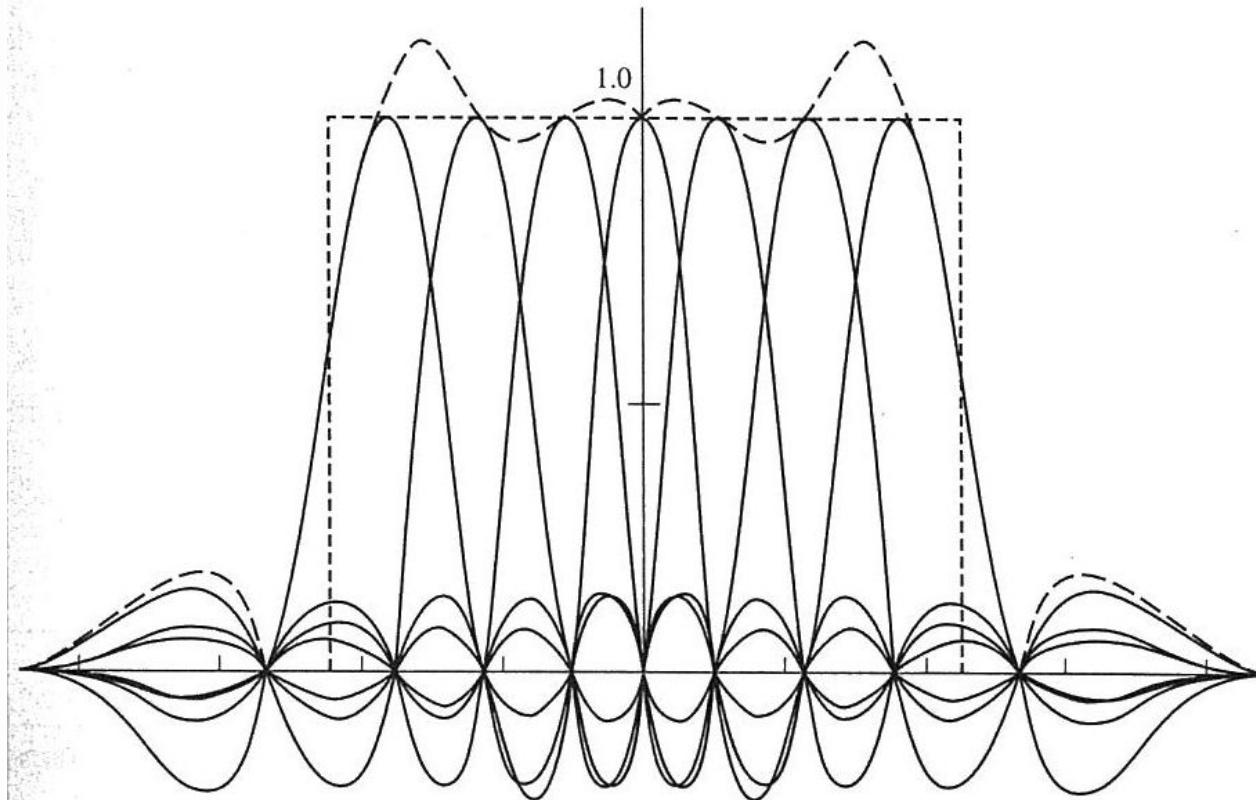
Since $l = 5\lambda$, $M = 5$ and $\Delta = 0.2$.

$$\theta_m = \cos^{-1}(m\Delta) = \cos^{-1}(0.2m), m = 0, \pm 1, \pm 2, \dots, \pm 5$$

m	θ_m	b_m	m	θ_m	b_m
0	90	1			
1	78.46	1	-1	101.54	1
2	66.42	1	-2	113.58	1
3	53.13	1	-3	126.87	1
4	36.87	0	-4	143.13	0
5	0	0	-5	180	0



Woodward-Lawson Method: Example (2)



- Desired pattern
- Line-source $|SF(\theta)|$ ($l = 5\lambda$)
- Composing functions $s_m(\theta)$,
 $m = 0, \pm 1, \pm 2, \pm 3$

Composing functions for line-source ($l = 5\lambda$)



Woodward-Lawson Method: Linear array

The pattern of each sample (uniform array) can be written as
(assuming $l = Nd$)

$$f_m(\theta) = b_m \frac{\sin\left[\frac{N}{2}kd(\cos\theta - \cos\theta_m)\right]}{N \sin\frac{kd}{2}(\cos\theta - \cos\theta_m)}$$

Composing function

For an odd number elements, the array factor becomes

$$\text{AF}(\theta) = \sum_{m=-M}^M b_m \frac{\sin\left[\frac{N}{2}kd(\cos\theta - \cos\theta_m)\right]}{N \sin\frac{kd}{2}(\cos\theta - \cos\theta_m)}$$

b_m can be obtained from the value at the sample points θ_m , i.e.,

$$b_m = \text{AF}(\theta = \theta_m)_d$$



Woodward-Lawson Method: Linear array (2)

The location of each sample is given by

$$\cos \theta_m = m\Delta = m\left(\frac{\lambda}{Nd}\right), \quad m = 0, \pm 1, \pm 2, \dots \text{ for odd samples}$$

$$\cos \theta_m = \begin{cases} \frac{2m-1}{2}\Delta = \frac{2m-1}{2}\left(\frac{\lambda}{Nd}\right), & m = 1, 2, \dots \\ \frac{2m+1}{2}\Delta = \frac{2m+1}{2}\left(\frac{\lambda}{Nd}\right), & m = -1, -2, \dots \end{cases} \text{ for even samples}$$

The normalized excitation coefficient of each element is given by

$$a_n(z') = \frac{1}{N} \sum_{m=-M}^M b_m e^{-j k z'_n \cos \theta_m}$$



Woodward-Lawson Method: Example

Same as Example 7.2; for $N=10$ and $d = \lambda/2$.

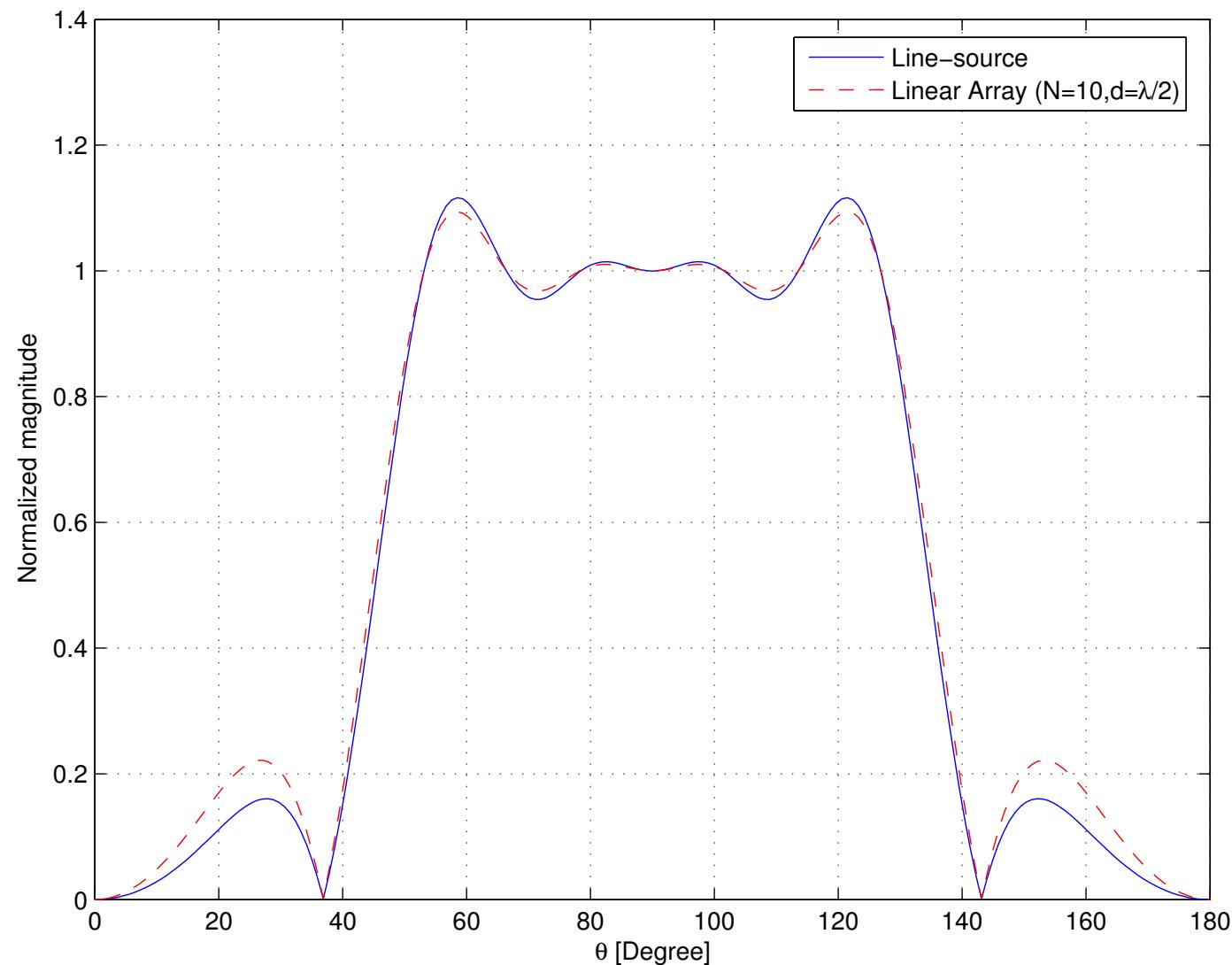
The coefficients can be found to be

Element number n	Position z'_n	Coefficient a_n
± 1	$\pm 0.25\lambda$	0.5696
± 2	$\pm 0.75\lambda$	-0.0345
± 3	$\pm 1.25\lambda$	-0.1001
± 4	$\pm 1.75\lambda$	0.1108
± 5	$\pm 2.25\lambda$	-0.0460

To obtain the normalized amplitude pattern of unity at $\theta=\pi/2$,
the array factor has been divided by $\sum a_n = 0.4998$

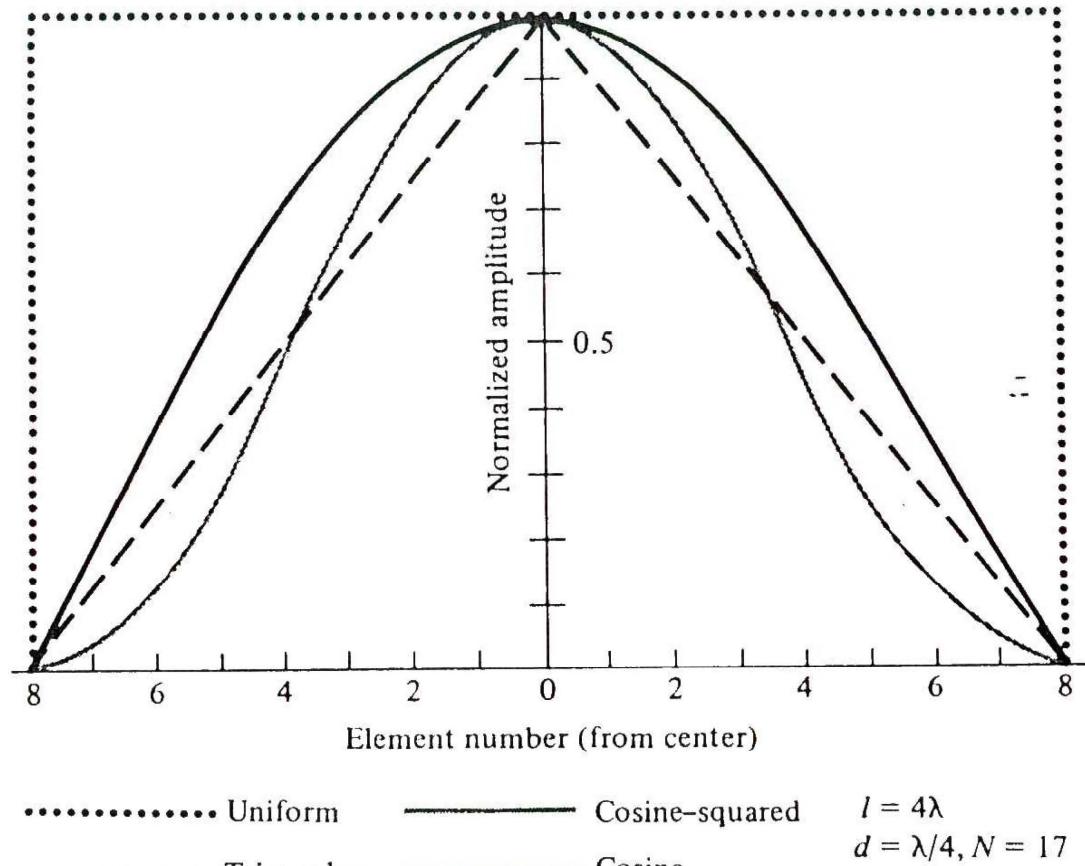


Woodward-Lawson Method: Example





Triangular, cosine, cosine squared distributions



(b) Amplitude distribution of uniform, triangular, cosine, and cosine squared discrete-element arrays

Fig 7.12b

Table 7.1 RADIATION CHARACTERISTICS FOR LINE SOURCES AND LINEAR ARRAYS
WITH UNIFORM, TRIANGULAR, COSINE, AND COSINE-SQUARED DISTRIBUTIONS

Distribution	Uniform	Triangular	Cosine	Cosine-Squared
Distribution I_n (analytical)	I_0	$I_1 \left(1 - \frac{2}{l} z' \right)$	$I_2 \cos \left(\frac{\pi}{l} z' \right)$	$I_3 \cos^2 \left(\frac{\pi}{l} z' \right)$
Distribution (graphical)				
Space factor (SF) $u = \left(\frac{\pi l}{\lambda}\right) \sin \theta$	$I_0 l \frac{\sin(u)}{u}$	$I_1 \frac{l}{2} \left[\frac{\sin\left(\frac{u}{2}\right)}{\frac{u}{2}} \right]^2$	$I_2 l \frac{\pi}{2} \frac{\cos(u)}{(\pi/2)^2 - u^2}$	$I_3 \frac{l}{2} \frac{\sin(u)}{u} \left[\frac{\pi^2}{\pi^2 - u^2} \right]$
Space factor SF				
Half-power beamwidth (degrees) $l \gg \lambda$	$\frac{50.6}{(l/\lambda)}$	$\frac{73.4}{(l/\lambda)}$	$\frac{68.8}{(l/\lambda)}$	$\frac{83.2}{(l/\lambda)}$
First null beamwidth (degrees) $l \gg \lambda$	$\frac{114.6}{(l/\lambda)}$	$\frac{229.2}{(l/\lambda)}$	$\frac{171.9}{(l/\lambda)}$	$\frac{229.2}{(l/\lambda)}$
First side lobe max. (to main max.) (dB)	-13.2	-26.4	-23.2	-31.5
Directivity factor (l large)	$2 \left(\frac{l}{\lambda}\right)$	$0.75 \left[2 \left(\frac{l}{\lambda}\right) \right]$	$0.810 \left[2 \left(\frac{l}{\lambda}\right) \right]$	$0.667 \left[2 \left(\frac{l}{\lambda}\right) \right]$

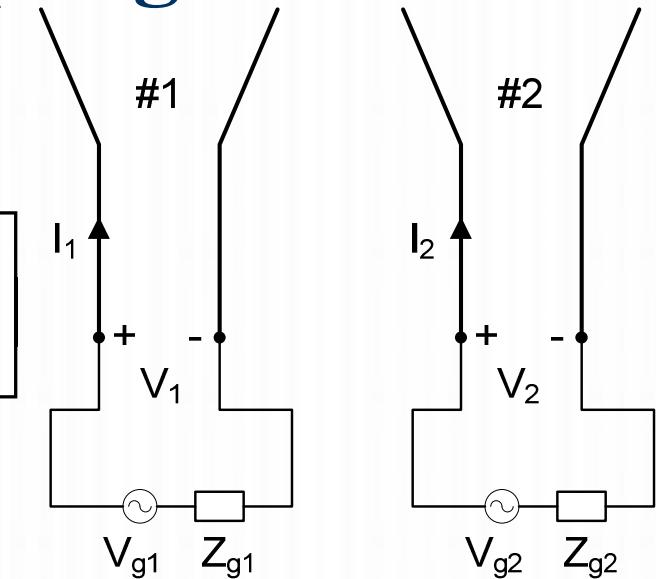


Mutual Coupling

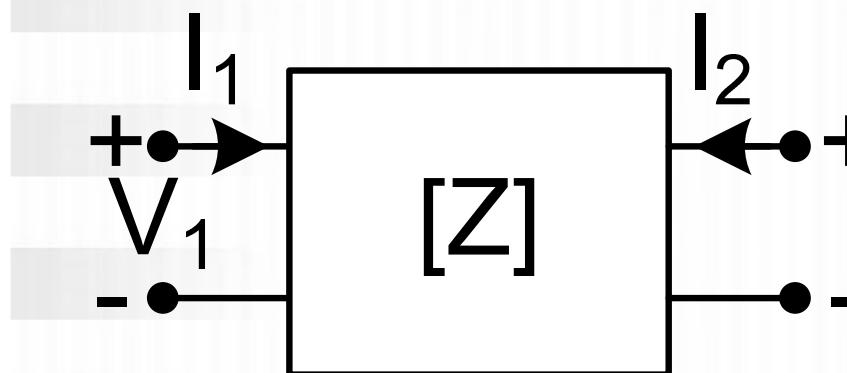
- Consider two antennas

$$\begin{aligned}V_1 &= Z_{11}I_1 + Z_{12}I_2 \\V_2 &= Z_{21}I_1 + Z_{22}I_2\end{aligned}$$

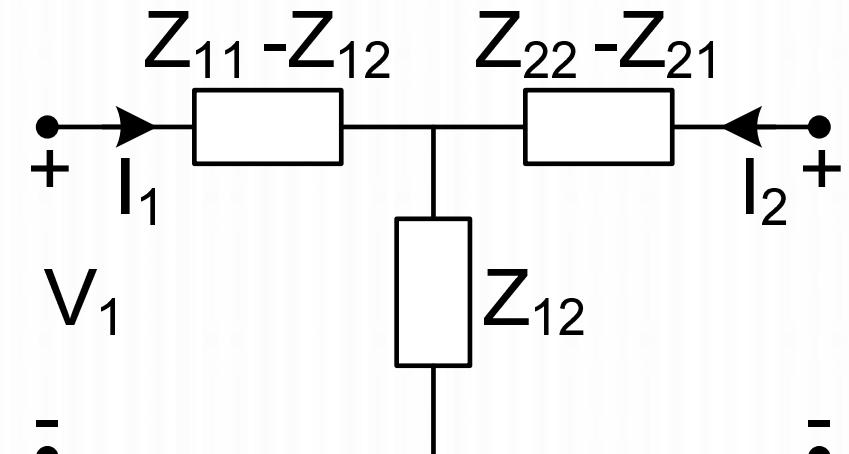
If reciprocal, $Z_{12} = Z_{21}$



Equivalent Circuit



Two-port Network



T-Network Equivalent Circuit



Mutual Coupling: 2 antennas

$$Z_{11} = \left. \frac{V_1}{I_1} \right|_{I_2=0} \quad Z_{12} = \left. \frac{V_1}{I_2} \right|_{I_1=0}$$

$$Z_{12} = Z_{21}$$

for reciprocal networks

$$Z_{21} = \left. \frac{V_2}{I_1} \right|_{I_2=0} \quad Z_{22} = \left. \frac{V_2}{I_2} \right|_{I_1=0}$$

$$Z_{11}, Z_{22} :$$

Input impedance

$$Z_{1d} = \begin{cases} \text{active} \\ \text{driving} \end{cases} \text{ point impedance}$$

$$Z_{1d} = \frac{V_1}{I_1} = Z_{11} + Z_{12} \frac{I_2}{I_1}$$

$$Z_{2d} = \frac{V_2}{I_2} = Z_{22} + Z_{21} \frac{I_1}{I_2}$$

Note: Z_{1d} depends on $\frac{I_2}{I_1}$; Z_{2d} depends on $\frac{I_1}{I_2}$



Mutual Coupling: 2 antennas (2)

- As I_1 and I_2 change, the driving point impedance changes.
- In a uniform array, the phase of I_1 and I_2 is changed to scan the beam.
- As the beam is scanned, the driving port impedance in each antenna changes.
- In general, $Z_{11}, Z_{12}=Z_{21}, Z_{22}$ can be calculated using numerical techniques.
- For some special cases, they can be calculated analytically.



Mutual Coupling: N antennas

$$\begin{bmatrix} V_1 \\ V_2 \\ \vdots \\ V_N \end{bmatrix} = \begin{bmatrix} Z_{11} & Z_{12} & \cdots & Z_{1N} \\ Z_{21} & Z_{22} & \cdots & Z_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ Z_{N1} & Z_{2N} & \cdots & Z_{NN} \end{bmatrix} \begin{bmatrix} I_1 \\ I_2 \\ \vdots \\ I_N \end{bmatrix} \quad \underline{\mathbf{V}} = \underline{\mathbf{Z}} \underline{\mathbf{I}}$$

$Z_{ij} = \frac{V_i}{I_j} \Big|_{I_k=0, k \neq j}$

In an N -element uniform array $I_n = I_0 e^{j(n-1)\beta}$

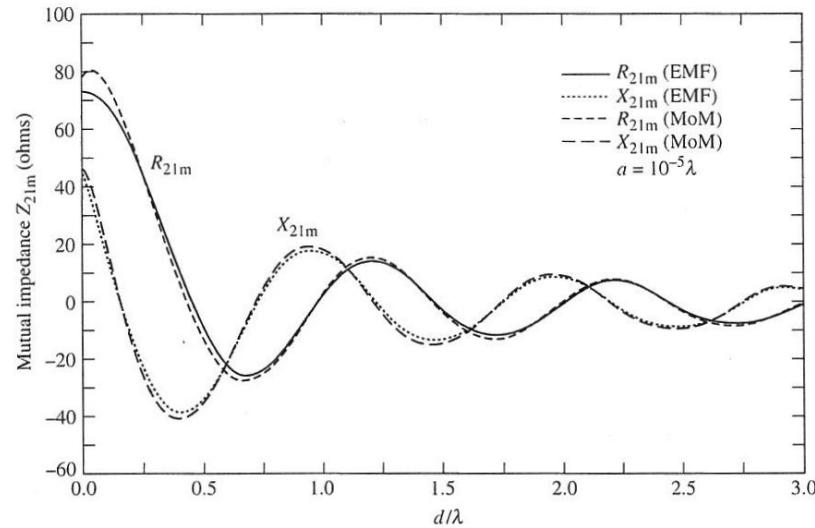
$$Z_{1d} = \frac{V_1}{I_1} = Z_{11} + Z_{12} e^{j\beta} + \cdots + Z_{1N} e^{j(N-1)\beta}$$

$$= \sum_{n=1}^N Z_{1n} e^{j(n-1)\beta}$$

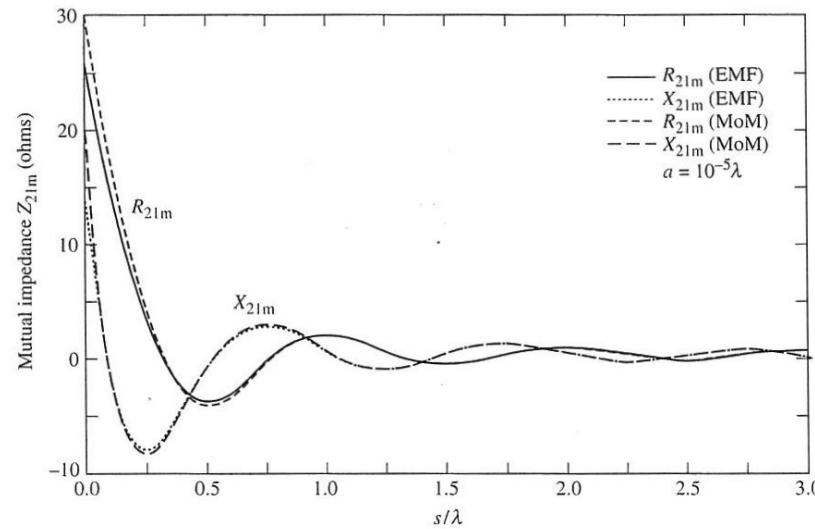
Note: Z_{1d} changes as β changes.



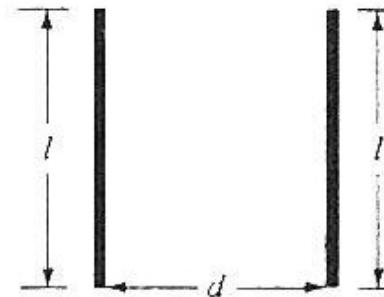
Mutual Coupling: 2 dipoles



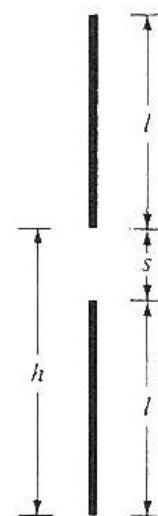
(a) Side-by-side



(b) Collinear



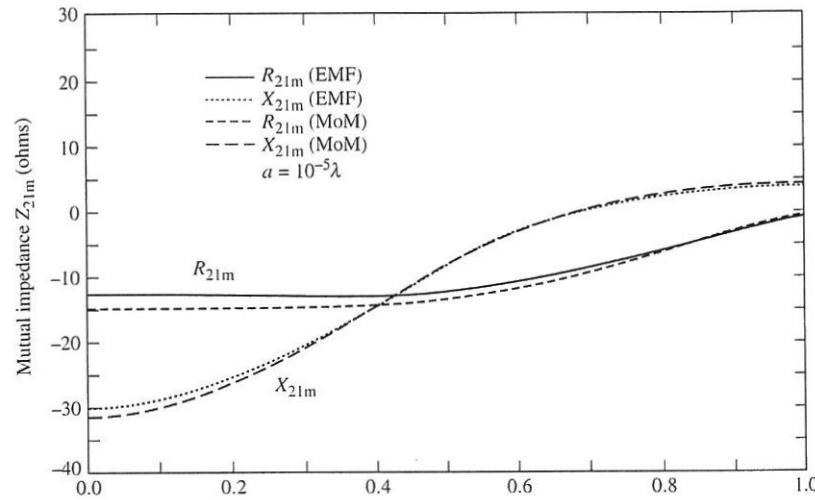
(a) Side-by-side



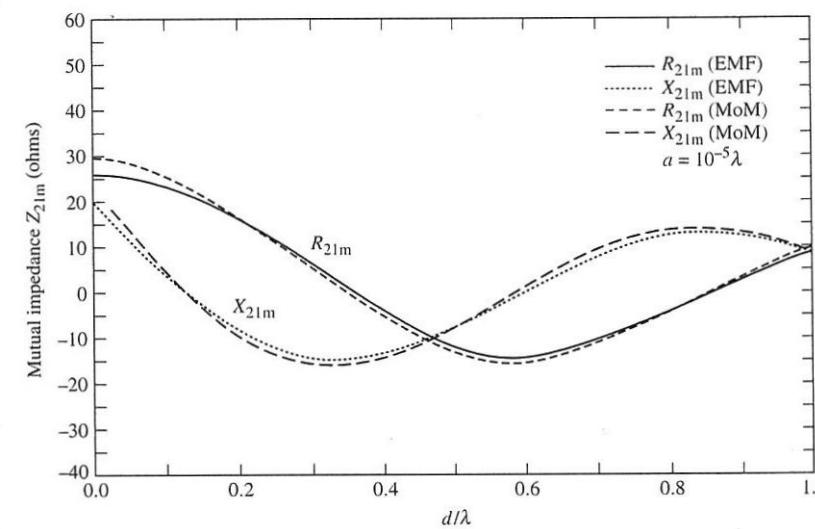
(b) Collinear



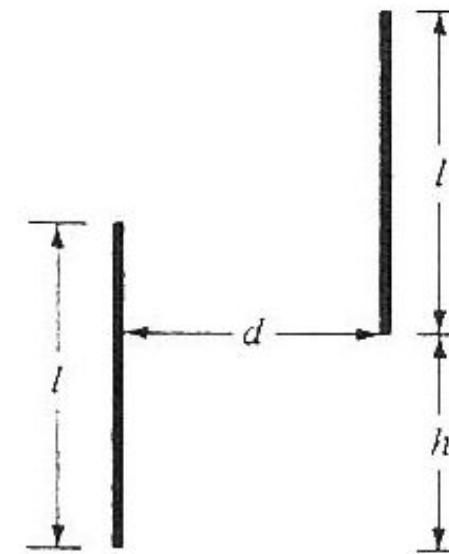
Mutual Coupling: 2 dipoles (2)



(a) $d = 0.5\lambda$



(b) $h = 0.5\lambda$



(c) Parallel-in-echelon